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Modeling the interaction of coolant flows at the liquid-solid boundary with allowance for the laminar boundary layer

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ABSTRACT

The forces acting on the elementary volume of liquid in the pipeline in the laminar boundary layer are considered. A new dimensionless complex the surface number and the concept of the turbulence coefficient laminar boundary layer is proposed. The calculation of Froude, Euler numbers, the inverse Reynolds number and the surface number in laminar boundary layer for water under normal conditions is given. It is shown that the Froude number and the inverse Reynolds number are 4-5 orders of magnitude smaller than the surface number and Euler number in laminar boundary layer, which allows neglecting the forces of gravitation and friction in these conditions. Equations are proposed for calculating the average thickness of laminar boundary layer. The dependence of the surface number on the coefficient of surface tension of the refrigerant is obtained. It is shown that a decrease in the surface tension coefficient minimizes the average thickness of the laminar boundary layers in the wall system of the pipeline and liquid and increases the average velocities of the coolant flows in these layers, as a result of which such a system is capable of more efficient transfer of heat. It is substantiated and experimentally confirmed that at optimum concentrations of surfactants, the values of the surface number are minimal.

Keywords: Average Thickness of the Laminar Boundary Layers, Surface Number, Turbulence Coefficient, Surfactants, Coefficient of Surface Tension.

1. INTRODUKTION

In the current energy crisis, the issue of energy efficiency of heat exchange processes and equipment for their provision in food, chemical, pharmaceutical, processing and other technologies becomes decisive. To increase the heat transfer coefficients in heat exchangers, various methods are used, in particular, the modification of the structural elements of boilers and other equipment [1], the increase in the turbulence of the refrigerant flows [2], the use of liquids with the optimal concentration of surfactants (SAS). For example, the maximum rate of heat exchange was observed when a nonionic surfactant was added to water [3]. The authors of [3] believed that the maximum rate of heat exchange in the first place may be due to the fact that this additive has a minimum capacity for the formation of foam. In addition, it is known that surfaceactive substances significantly, approximately 2 times, reduce the coefficients of surface tension of water and other liquids.

2. PROBLEM FORMULATION

After analyzing the literature sources, we came to the conclusion that the rate of heat exchange in liquid refrigerants through the laminar boundary layer (LBL) depends on the following main factors:

- laminar flow is responsible for the turbulent flow, but is less energy-efficient [2,4];

- LBL, namely its average thickness is responsible for the total thermal resistance of the system [5,6];

- the thermal resistance depends on the coefficient of surface tension of the refrigerant [7];

- the intensity of heat exchange depends on the hydrophilicity or hydrophobicity of the wetting surface [8];

Based on the foregoing, the purpose of this article was to offer a model for the interaction of refrigerants with a separating solid wall, which will cover all the above factors as much as possible.

3.1. Analysis of the forces in LBL

It is common knowledge that for the pipeline, and in particular for tubes heat exchangers the vectors of flow average velocities are distributed in tube longitudinal section as a parabola (Fig. 1).

At the boundary of the flow of a liquid and a pipe there is always a LBL [9]. This layer has a very small average thickness, but its effect on heat transfer and diffusion processes that occur in the flow is crucial. Consider the elementary volume of liquid in the pipeline within the limits LBL (Fig. 1.)



Figure 1. Forces acting on an elementary volume of liquid in LBL

On elementary volume of liquid in LBL is acted upon by forces:

- 1. Force of surface tension of liquid: $F_{\sigma} = 2\pi (dx) \cos \theta \sigma$;(1)
- 2. Force of gravity: $mg = \rho g(dxdydz)$; (2)

3. Friction force:
$$T = \mu \frac{d^2 V_z}{dz^2} (dx dy dz);$$
 (3)

4. Force of inertia:
$$F_i = ma = \rho \frac{dV_z}{d\tau} (dxdydz);$$
 (4)

5. Force of pressure:
$$P = -\frac{d\mathbf{p}}{dz}(dxdydz)$$
, (5)

where, P – fluid pressure acting on the upper face of the elementary volume, Pa; V_z – average fluid velocity in LBL, m/s; σ – the surface tension of the liquid, N/m; ρ – fluid density, kg/m³; g – acceleration of gravity, m/s²; μ – coefficient of dynamic viscosity, N.s/m²; τ – time, s; $cos\theta$ – cosine of wetting angle;

According to the principle of d'Alembert, the algebraic sum of all forces acting on the elementary volume is equal to the force of inertia. When dividing by (dxdydz), we get the equation 6. This is the Navier-Stokes equation, which is supplemented by surface forces, which according to our statement in LBL reach commensurate values with pressure forces:

$$\frac{2\pi (dx)\sigma\cos\theta}{(dxdydz)} + \rho g - \frac{dP}{dz} + \mu \frac{d^2 V_z}{dz^2} = \rho \frac{dV_z}{d\tau}$$
(6)

Equation 6 can not be integrated, so we obtain a number equation from it, applying the similarity theory. The symbols of differentiation of the differential equation and direction are removed, we replace the linear parameters of the elementary volume dx, dy, dz by l.

Based on the similarity theory, taking into account that the fluid velocity and the linear parameter *l* inside LBL are very small, dividing the right and left sides of equation (6) by ${}^{\rho V_z}/_{\tau}$, we get the numbers: we get the numbers:

1.
$$\frac{2\pi\sigma\cos\theta\tau}{\rho V_z l^2} = \frac{2\pi\sigma\cos\theta\mu}{\rho V_z^2 l\mu} = \frac{1}{N}\frac{2\pi\sigma\cos\theta}{\mu V_z} = Po-\text{ supface}$$

number in LBL (7)

where N — modified Reynolds number in LBL $N = \frac{\rho V_z l}{\mu} \approx 10,47 - 11,5$ [10];

2.
$$\frac{p\tau}{\rho V_z l} = \frac{\Delta p}{\rho V_z^2} = Eu$$
 – Euler number in LBL (8)

3.
$$\frac{\rho g \tau}{\rho V_z} = \frac{g \tau l}{V_z l} = \frac{g l}{V_z^2} = Fr$$
 – Froude number in LBL (9)

4.
$$\frac{\mu V_z \tau}{l^2 \rho V_z} = \frac{\mu}{l \rho V_z} = \frac{1}{\text{Re}}$$
 – Reynolds number in LBL (10)

We substitute in the numbers 1, 2, 3, 4 characteristics of water under normal conditions: $\Delta P \approx 100000$ Pa – decline in pressure in the pipeline; $\rho \approx 1000$ fluid density, kg/m³; $V \approx 1$ m/s – the average velocity of liquids in pipelines; $V_z \approx \frac{V}{10} \approx 0.1 [10]$ m/s – the average velocity of liquids in LBL; $\sigma \approx 0.0725$ N/m – the surface tension of the water under normal conditions, N/m; $\cos \theta \approx 0.8$ – cosine of wetting angle; $\mu \approx 1.10^{-3}$ – coefficient of dynamic viscosity, N.s/m²; $\delta \approx 0.3 \cdot 10^{-3}$ m – average thickness LBL under turbulent driving

1.
$$P_{O} = \frac{1}{N} \frac{2\pi \cdot \sigma \cdot \cos\theta}{\mu V_{z}} \approx \frac{1}{10.5} \frac{2 \cdot 3.14 \cdot 0.0725 \cdot 0.8}{1 \cdot 10^{3} \cdot 0.1} \approx 347;$$

2. $Eu = \frac{\Delta p}{\rho V_{z}^{2}} \approx \frac{100000}{1000 \cdot (0.1)^{2}} \approx 10000;$
3. $Fr = \frac{g \cdot l}{V_{z}^{2}} \approx \frac{10 \cdot \delta}{(0.1)^{2}} \approx \frac{10 \cdot 0.3 \cdot 10^{-3}}{(0.1)^{2}} \approx 0.3;$
4. $\frac{1}{Re} = \frac{\mu}{\delta \rho V} \approx \frac{10^{-3}}{0.3 \cdot 10^{-3} \cdot 1000 \cdot 0.1} \approx 0.03;$

condition of water in the pipes.

The Froude number and the inverse Reynolds number are 4-5 orders of magnitude smaller than the surface criterion and the Euler number in the laminar boundary layer, which allows neglecting the forces of gravity and friction in these conditions.

3.2 The formula for calculating the average thickness of LBL

Consider the pressure force and force of surface tension in laminar boundary layer actings in the elementary ring of liquid (Fig. 1):

- the area of elementary ring: $dA = 2\pi r dr$;

- the pressure force acting on elementary ring of the laminar boundary layer: $dP = \Delta P 2\pi r dr$;

- the force of surface tension acting on elementary ring in the laminar boundary layer: $F_{\sigma} = 2\pi\sigma\cos\theta dr$;

Taking into account only the pressure forces and the forces of surface tension, we integrate the equation changing the limits of integration r and δ .

The physical meaning of integration in this system is that the laminar boundary layer covers the entire passage of the pipeline, the boundary of the transition laminar-turbulent, have the limiting case-the average thickness of the laminar boundary layer is equal to the radius of the tubes $\delta = r = \frac{d}{2}$ (Fig. 1).

$$P = \Delta P 2\pi \int_{0}^{\delta} r dr = \pi \Delta P \delta^{2};$$

$$dF_{\sigma} = 2\pi\sigma \cos\theta \int_{0}^{r} d\delta = 2\pi\sigma \cos\theta r;$$

Based on the principle of theoretical fluid mechanics, after integration and reduction of the number π , we get equation (11):

$$\delta = \sqrt{\frac{\sigma d \cos \theta}{\Delta P}} \,. \tag{11}$$

Formula (11) describes the flow of liquid at the boundary of the laminar and transient flow of liquid flow at a critical Reynolds number $Re_{cr} = 2320$. For Reynolds number values that are greater than the critical values in equation (11) introduced the coefficient of turbulence K_T (12).

$$\delta = \frac{\sqrt{\frac{\sigma \cos\theta d}{\Delta P}}}{K_T} , \qquad (12)$$

 $K_T = \frac{Re}{Re_{cr}}$

One should notice that for the boundary conditions, when the pipe is reduced to the capillary size, i.e. the diameter *d* becoming small, the equation (12) transforms into the Laplace's equation $\Delta P = \frac{4\sigma\cos\theta}{d}$; (at $\delta = r = \frac{d}{2}$; $d^2 = \sigma d\cos\theta$

 $\frac{d^2}{4} = \frac{\sigma d \cos \theta}{\Delta P}$); It is obvious that in the capillary goes

laminar fluid motion mode. Therefore, the ratio $\frac{Re}{Re_{cr}} = 1$.

If write the difference of pressure in accordance with Darcy's equation and make conversions, from equation (12) we obtain the equation (13):

$$\delta = \frac{\frac{d^2}{\text{Re}\,\mu}\sqrt{\frac{2\sigma\cos\theta\rho}{fL}}}{\frac{\text{Re}}{\text{Re}_{cr}}} \tag{13}$$

where ΔP – pressure drop along the pipe or apparatus, Pa; r – the radius of elementary ring, m; f – Darcy coefficient; ρ – fluid density, kg/m³; L – length of the pipe, m;

d – diameter of the pipe, m; $cos\theta$ – the cosine of the contact angle; δ – the average thickness of the laminar boundary layer, m; σ – the surface tension of the liquid, N/m; μ – coefficient of dynamic viscosity, N.s/m²; Re – the Reynolds number; Re_{cr} – the critical Reynolds number (Re_{cr} , ≈ 2320);

V- the velocity of the fluid flow, m/s.

3.3 Effect of surfactants on the average thickness of LBL

The resulted ratios (12, 13) shows that the average thickness LBL marginally depends on the coefficient of fluid dynamic viscosity. And it is not unexpected, because the friction forces in laminar boundary layer are very small due to the very lows sliding velocity of adjacent layers. However, the average thickness the laminar boundary layer depends on the surface tension coefficient and of the cosine of the contact.

Equation (7) shows that the numerical values of the surface criterion depend substantially on the surface tension coefficient. It is known that the coefficient of surface tension of liquid-phase heat carriers can be reduced by introducing optimal concentrations of SAS. To the aqueous solutions (model fluid 1) we added the most common non-ionic SAS. As a surfactant for the components of milk (model liquid 2) - vegetable oil of pumpkin. To determine the range of values of the surface criterion, a series of experiments was performed to ensure changes in the coefficient of surface tension, the cosine of the wetting angle, and the dynamic coefficient of viscosity of aqueous solutions and milk under the influence of SAS, which were measured according to well-known procedures.

In Fig. 2 shows the dependence of the surface tension coefficient, the dynamic coefficient of viscosity of the water (a) and milk (b) wetting angle on the SAS concentrations. As can be seen from the graphs, the curve of the dependence of the surface tension on the concentration has minima close to the critical concentration of micelle (CCM) formation. That is, with a slight increase in surfactant concentrations, the surface tension coefficient decreases sharply to the CCM, and at a concentration above the CCM, its decrease is insignificant. For non-ionic SAS the CCM is observed at a concentration $(0.05 \dots 0.10)$ by mass. %. At the same time, the coefficient of surface tension decreases by 2.32 times in comparison with water.

As can be seen from the graphs, the minimum value of the surface tension coefficient of milk is observed at a concentration (0.5...0.6) by mass. % pumpkin oil. These concentrations SAS were considered optimal. At these SAS concentration, the value for the dynamic viscosity coefficient was also selected.

Let us show the change in the surface number in LBL for aqueous solutions with the addition of the optimum concentrations of the surfactants studied. The rate in LBL was determined from the modified Reynolds number, taking the value N = 10.5, and the average thickness of the LBL was found from the formula (13). In the food, pharmaceutical and processing industries, the average fluid velocity in heat

exchange equipment is $V \approx 1 m/s$, pipe diameter $d = 21 \cdot 10^{-3}$ m, pipe length L=3 m.



Figure 2. Dependence of the coefficient of surface tension and the dynamic coefficient of viscosity on the concentration of surfactants (a) water from the concentration of non-ionic SAS, (b) milk from the concentration of pumpkin oil.

The results of calculations are presented in Table 1.

Let us consider the cross section of the flow in a pipeline under the turbulent (T) regime of fluid motion (Fig. 3). The velocity vectors in it are distributed like a parabola, but with a wider vertex.

Reducing the average thickness LBL of water and milk, the speed in them increases, and this intensifies the passage of the amount of heat through it (Fig. 4).

 Table 1. The value of the surface number at the optimum

 surfactant concentrations in the model fluid 1 ("ice" water)

 and in the model fluid 2 (milk)

Coolant	δ , the	V_z , the	P_o , the
	average	average	surface
	thickness	speed	number
	LBL, mkm	in LBL,	in LBL
		m/s	
Water	116	0.116	254
Water + $(0.05 \dots 0.10)$	79	0.167	86
mass. %. non-ionic			
SAS			
Milk	113	0.141	128
Milk + (0,5 0,6)	82	0.155	106
mass. % pumpkin oil			



Figure 3. Medium velocity vectors in LBL without SAS and with using SAS



Figure 4. Heat transfer scheme Q through the metal wall from milk to water

4. CONCLUSIONS

1. LBL has a powerful field of surface tension forces in the liquid.

2. The notion of a surface number in LBL is introduced. The intensity of surface forces is characterized by the magnitude of the surface number.

3. The proposed formulas 11, 12, 13 for calculations of the average thickness of LBL and modeling the interaction of coolant flows at the liquid-solid boundary with allowance for the laminar boundary layer are, in our opinion, acceptable, since they simultaneously cover all the factors listed above:

- LBL, namely its average thickness is responsible for the total thermal resistance of the system;

- the thermal resistance depends on the coefficient of surface tension of the refrigerant;

- the intensity of heat exchange depends on the hydrophilicity or hydrophobicity of the wetting surface;, dependencies on which were obtained by different authors at different times.

4. Formulas 11, 12, and 13, in our opinion, are reliable, since under boundary conditions they are transformed into the classical Laplace relation.

5. The use of appropriate surfactants of the optimum concentrations to liquid heat carriers will significantly increase the overall heat transfer coefficient in the heat exchange equipment of food, pharmaceutical, processing and other industries.

6. The optimal concentrations of surfactants for two model fluids (milk and water) were experimentally found, at which the average thicknesses of LBL is minimal.

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NOMENCLATURE

A	area of elementary ring, m ²
$\cos\theta$	cosine of wetting angle
d	diameter of the pipe, m
Eu	Euler number in LBL
Fr	Froude number in LBL
F_{σ}	surface tension force of the liquid, N
F_i	force of inertia, N
ſ	hydraulic frictiocn coefficient
Ĺ	length of the pipe, m
N	modified Reynolds number in LBL
Р	force of pressure, N
р	fluid pressure acting on the upper face of
	the elementary volume, Pa
ΔP	pressure drop along the pipe or unit, Pa
P_o	supface number in LBL
Re	Reynolds number in LBL
Т	friction force, N
V	velocity, m/s
V_{z}	average fluid velocity in LBL, m/s
*	

Greek symbols

δ	average thickness of the laminar boundary
	layer, m
μ	dynamic viscosity coefficient, N.s/m ²
ρ	fluid density, kg/m ³
σ	surface tension of the liquid, N/m
τ	time, s

Subscripts

aver	average
Cr	critical

Abbreviations

CCM	critical concentration of micelle
LBL	laminar boundary layer
SAS	surfactants