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Variable Temperature and Concentration Impacts on Radiative Chemically Magnetohydrodynamic Viscoelastic Fluid Flow Through Porous Moving Plate



Karna Suneetha¹, Shaik Mohammed Ibrahim^{2*}, Gurram Venkata Ramana Reddy¹, Prathi Vijaya Kumar²

¹ Department of Mathematics, K. L. E. F, Vaddeswaram, Guntur, Andhra Pradesh 522502, India ² Department of Applied Mathematics, GITAM Deemed to be University, Andhra Pradesh 530045, India

Corresponding Author Email: sibrahim@gitam.edu

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ABSTRACT

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The flow has been made by considering variable temperature and radiation effects for the magnetohydrodynamic viscoelastic fluid past a moving vertical plate in a porous medium. Chemical reaction and concentration have been taken into account. The governed mathematical statement is handled analytically by perturbation technique. The main view of this research is to investigate the effects of parameters and numbers in the problem on fluid flow, thermal boundary and concentration profiles. The velocity profile has been reduced by increasing the magnetic parameter due to the Lorentz force in the opposite direction of flow. Temperature profile is increased by rising thermal radiation and concentration distribution is decreased by enhancing the chemical reaction and Schmidt number. The Schmidt number represents the relative ease of the molecular momentum and mass transfer and it is very important in multiphase flows. The effect of increasing values of the Schmidt number is to reduce the momentum boundary layer and this leads to the thinning of the diffusion layer. Furthermore, at the end of this paper the effects of different parameters on skin friction coefficient and local Nusselt number are investigated.

1. INTRODUCTION

Convection is of fundamental interest in numerous engineering, industrial, and environmental applications such as cooling of electronic devices, air-conditioning systems, atmospheric flows, and security of energy systems and in designs related to thermal insulation. Flow, of combined thermal and mass transfer in a porous media has several industrial applications such as filtration process and powerengineering equipment such as cooling of electronic devices, microelectronic chips, printed circuit boards and photovoltaic sheets. It is also important in various engineering and geophysical problems. In numerous engineering and technological applications, importance of non-Newtonian fluids cannot be negated. Examples of these materials may include shampoos, mayonnaise, blood, paints, alcoholic beverages, yogurt, cosmetics, and syrups etc. Mathematical modelling of these fluids is very tedious as typical Navier-Stokes equations are not enough to express characteristics of non-Newtonian fluids. These fluids are categories as differential, rate and integral types. Viscoelastic fluids are subclasses of non-Newtonian fluids which possess memory effect. These fluids exhibit certain amount of energy which is responsible for the partial elastic recovery upon the removal of stress. Beard and Walters [1] first perceived the boundary layer analysis of idealized visco-elastic fluid. Natural convection flow between two vertical parallel plates was proposed by Singh et al. [2]. Sajid et al. [3] developed a fully mixed convection flow between two permeable vertical walls in visco-elastic. The importance of visco-elastic fluid flow in presence of different parameters have studied by many authors [4-14]. Radiative dissipative MHD natural convection flow under the influence of heat source and sink was derived by Suneetha et al. [15].

The effect of thermal radiation becomes significant for several industrial processes such as glass production, furnace design, electrical power generation and solar power technology. A good working knowledge of thermal radiation helps in designing of important equipments such as design of fins, ceramic and glass producing units and various propulsion devices for aircraft, missiles, satellites, space vehicles etc. Keeping in mind such importance, Chemical reaction is notable in several procedures like chemical processing, hydrometallurgical industry, fibrous insulation, atmospheric flows, damage of crops because of freezing, water and air pollutions, production of ceramics and polymer, fog formation and dispersion etc. Reddy et al. [16] studied the proposed the magnetohydrodynamic free convection flow behaviour in a porous medium with constant heat and mass flux under thermal radiation and chemical reaction. Ahmed [17] and Sandeep et al. [18] observed the nature of chemically reactive flow over a vertical plate under different background.

In certain fluid applications, working fluid heat source or sink effects are important. Sample studies dealing with these effects have been reported by many authors. In recent past, great care has been taken to audit the repercussion of chemical reaction and heat source (or sink) on different flow types [19-23]. The grasp on this subject assist to reconcile abundant biological problems. Considering the model of visco-elastic fluid, many scientists have solved problems of engineering interests viz. In the last few years, many investigations [24-38], have been carried out regarding the present work. In the light of the above studies, the objective of the current investigation is to prove the influence of heat generation or absorption and first order chemical reaction effects on laminar boundary layer flow through porous medium with thermal radiation, variable temperature and concentration. The dimensionless equations are then solved analytically using perturbation technique. The behaviors of different parameters on the physical quantities have been examined.

2. GOVERNING EQUATIONS

A two-dimensional unsteady MHD flow of an incompressible electrically conducting fluid over a semiinfinite vertical permeable moving plate permeable stretching surface in presence of thermal radiation is considered. The system of coordinate is taken in way that *x*-axis is measured along the sheet and y-axis is orthogonal to it as presented in Figure 1. Induced magnetic field is negligible as compared to the applied magnetic field. We assume that the equations are subjected to visco-elastic fluid flow proposed by Babu et al. [39]. In the absence of the gradient of pressure, the governing equations expressing conservation of mass, momentum, energy and species are given as follows:



Figure 1. Physical model of the problem

$$\frac{\partial v^*}{\partial y^*} = 0 \tag{1}$$

$$\frac{\partial u^*}{\partial y^*} + v^* \frac{\partial u^*}{\partial y^*} = v \frac{\partial^2 u^*}{\partial y^{*2}} - \left(\frac{\sigma B_0^2}{\rho} + \frac{v}{K^*}\right) u^* + g \beta_T (T^* - T_\infty) + g \beta_C (C^* - C_\infty) - k_0 \left(\frac{\partial^3 u^*}{\partial t^* \partial y^{*2}} + v^* \frac{\partial^3 u^*}{\partial y^{*3}}\right)$$
(2)

$$\frac{\partial T^*}{\partial t^*} + v^* \frac{\partial T^*}{\partial y^*} = \frac{k}{\rho C_p} \frac{\partial^2 T^*}{\partial y^{*^2}} - \frac{1}{\rho C_p} \frac{\partial q_r^*}{\partial y^*} - \frac{Q_0}{\rho C_p} \left(T^* - T_\infty\right)$$
(3)

$$\frac{\partial C^*}{\partial t^*} + v^* \frac{\partial C^*}{\partial y^*} = D \frac{\partial^2 C^*}{\partial y^{*^2}} - k_1 \left(C^* - C_\infty \right) \tag{4}$$

The boundary conditions for the above described model

$$u^{*} = u_{p^{*}}, \ T^{*} = T_{w} + \varepsilon \left(T_{w} - T_{\infty}\right) e^{n^{*} t^{*}},$$
 (5)

$$C^* = C_w + \varepsilon \left(C_w - C_\infty \right) e^{n t^*}, \text{ at } y^* = 0$$

$$u^* = 0, \ T^* \to T_\infty, \ C^* \to C_\infty \text{ at } y^* \to \infty$$

It is unambiguous that Eq. (1) that the velocity of suction at the surface plate is time function. Presuming it yields into the form:

$$v^{*} = -V_{0}(1 + \varepsilon A e^{n^{*} t^{*}})$$
(6)

 ε and A are small such that $\varepsilon \ll 1, A \ll 1$. Acknowledging a self-similar solution of the form

$$u = \frac{u^{*}}{V_{0}}, u = \frac{v^{*}}{V_{0}}, y = \frac{V_{0}y^{*}}{v}, t = \frac{V_{0}^{2}y^{*}}{v}, u_{p} = \frac{u_{p}^{*}}{V_{0}},$$

$$n = \frac{n^{*}v}{V_{0}^{2}}, \theta = \frac{T^{*} - T_{\infty}}{T_{w} - T_{\infty}}, C = \frac{C^{*} - C_{\infty}}{C_{w} - C_{\infty}}$$
(7)

the basic field Eqns. (2) to (4) can be expressed in nondimensional form as

$$\frac{\partial u}{\partial t} - \left(1 + \varepsilon A e^{nt}\right) \frac{\partial u}{\partial y} = \frac{\partial^2 u}{\partial y^2} - \left(M + \frac{1}{K}\right) u + Gr\theta + GmC - E \left[\frac{\partial^3 u}{\partial t \partial y^2} - \left(1 + \varepsilon A e^{nt}\right) \frac{\partial^3 u}{\partial y^3}\right]$$
(8)

$$\frac{\partial\theta}{\partial t} - \left(1 + \varepsilon A e^{nt}\right) \frac{\partial\theta}{\partial y} = \frac{1}{\Pr} \frac{\partial^2 \theta}{\partial y^2} - \left(Q + R\right)\theta \tag{9}$$

$$\frac{\partial C}{\partial t} - \left(1 + \varepsilon A e^{nt}\right) \frac{\partial C}{\partial y} = \frac{1}{Sc} \frac{\partial^2 C}{\partial y^2} - KrC$$
(10)

$$u = u_p, \ \theta = 1 + \varepsilon e^{nt}, \ C = 1 + \varepsilon e^{nt} \text{ at } y = 0$$

$$u \to 0, \ \theta \to 0, \ C \to 0 \text{ as } y \to \infty$$
 (11)

$$Gr = \frac{(T_w - T_w)\beta_T g\upsilon}{V_0^3}, Gm = \frac{(C_w - C_w)\beta_C g\upsilon}{V_0^3},$$

$$R = \frac{4\upsilon}{\rho C_p V_0^2}, \Pr = \frac{\rho C_p \upsilon}{k}, K = \frac{K^* V_0^2}{\upsilon^2}, Kr = \frac{K_1 \upsilon}{V_0^2}, \quad (12)$$

$$Sc = \frac{\upsilon}{D}, Q = \frac{\upsilon Q_0}{\rho C_p V_0^2}, E = \frac{k_0 V_0^2}{\upsilon^2}$$

3. PROBLEM SOLUTION

Solutions of Eqns. (8) to (10) are reaped by regular and multi-parameter perturbation technique. E, ε and A are presumed small, such that $E \ll 1$ and $\varepsilon \ll 1$.

For getting solutions we introduce

$$u(y,t) = u_0(y) + \varepsilon e^{nt} u_1(y) + 0(\varepsilon^2)$$

$$\theta(y,t) = \theta_0(y) + \varepsilon e^{nt} \theta_1(y) + 0(\varepsilon^2)$$

$$C(y,t) = C_0(y) + \varepsilon e^{nt} C_1(y) + 0(\varepsilon^2)$$
(13)

where, u_{0} is the mean velocity, θ_0 is the mean temperature and C_0 is the mean concentration. Applying Eq. (13) into Eqns. (8) to (10). Tallying non-harmonic and harmonic statement to above location, after neglecting coefficient of ε^2 , we secure zero order

$$Eu_0''' + u_0'' + u_0' - n_1 u_0 = -Gr\theta_0 - GmC_0$$
(14)

$$\theta_0'' + \Pr \theta_0' - n_3 \Pr \theta_0 = 0 \tag{15}$$

$$C_o'' + ScC_0' - ScKrC_0 = 0 \tag{16}$$

where, $n_1 = M + \frac{1}{K}$, $n_3 = Q + R$ With

With

$$u_0 = u_p, \theta_0 = 1, C_0 = 1, \qquad \text{at } y = 0$$

$$u_0 \to 0, \theta_0 \to, C_0 \to 0, \qquad \text{as } y \to \infty \qquad (17)$$

And first order

$$Eu_{1}^{'''} + (1 - nE)u_{1}^{''} + u_{1}^{'} - n_{2}u_{1} = -Gr\theta_{1} - GmC_{1}$$

-AEu_{0}^{'''} - Au_{0}^{'} (18)

$$\theta_1'' + \Pr \theta_1' - n_4 \Pr \theta_1 = -\Pr A \theta_0' \tag{19}$$

$$C_1'' + ScC_1' - Scn_5C_1 = -AScC_0'$$
(20)

where,
$$n_2 = \left(M + \frac{1}{K} + n\right)$$
, $n_4 = Q + R + n$, $n_5 = Kr + n$.

With corresponding boundary conditions

$$u_1 = 0, \ \theta_1 = 1, \ C_1 = 1 \qquad \text{at } y = 0$$

$$u_1 \to 0, \theta_1 \to, C_1 \to 0, \qquad \text{as } y \to \infty \qquad (21)$$

Eqns. (14) and (18) are differential equations of 3rd order by virtue of visco-elastic parameter. Since there are exclusively two accessible boundary conditions, it necessitates an additional boundary condition to novel solution.

$$u_0(y) = u_{00}(y) + Eu_{01}(y) + 0(E^2)$$

$$u_1(y) = u_{10}(y) + Eu_{11}(y) + 0(E^2)$$
(22)

Put Eq. (22) in Eq. (14). Now compare the coefficient of first and zero order of E, Ne can procure

$$u_{00}'' + u_{00}' - n_1 u_{00} = -Gr\theta_0 - G_m C_0$$
⁽²³⁾

$$u_{01}'' + u_{01}' - n_1 u_{01} = -u_{00}'''$$
(24)

The boundary conditions are

$$u_{00} = u_p, u_{01} = 0, \text{ on } y = 0$$

 $u_{00} \to 0, u_{01} \to 0, \text{ as } y \to \infty$ (25)

Put Eq. (22) in Eq. (18) and compare the coefficients of zero and first order of E,

We get

$$u_{10}'' + u_{10}' - n_2 u_{10} = -Gr\theta_{01} - G_m C_1 - Au_{00}'$$
(26)

$$u_{11}'' + u_{11}' - n_2 u_{11} = -Au_{00}'' - Au_{01}' - u_{10}''' + nu_{10}''$$
(27)

with

$$u_{10} = 0, u_{11} = 0$$
 on $y = 0$
 $u_{10} \to 0, u_{11} \to 0$ as $y \to \infty$ (28)

Using the Eqns. (25) and (28) one can solve the Eqns. (23), (24), (26), and (27) in order to obtain

$$u_{00} = A_{7}e^{-m_{5}y} - A_{5}e^{-m_{3}y} - A_{6}e^{-m_{1}y}$$

$$u_{01} = A_{11}e^{-m_{6}y} + A_{8}e^{-m_{5}y} - A_{9}e^{-m_{3}y} - A_{10}e^{-m_{1}y}$$

$$u_{10} = A_{17}e^{-m_{7}y} + A_{12}e^{-m_{5}y} - A_{13}e^{-m_{4}y} - A_{14}e^{-m_{3}y}$$

$$-A_{15}e^{-m_{2}y} - A_{16}e^{-m_{1}y}$$

$$u_{11} = A_{25}e^{-m_{8}y} + A_{18}e^{-m_{7}y} + A_{19}e^{-m_{6}y} + A_{20}e^{-m_{5}y} - A_{21}e^{-m_{4}y}$$

$$-A_{22}e^{-m_{3}y} - A_{23}e^{-m_{2}y} - A_{24}e^{-m_{1}y}$$

$$u_{0}(y) = u_{00}(y) + Eu_{01}(y)$$

$$u_{0}(y) = \left(A_{7}e^{-m_{5}y} - A_{5}e^{-m_{3}y} - A_{6}e^{-m_{1}y}\right)$$

$$+E\left(A_{11}e^{-m_{6}y} + A_{8}e^{-m_{5}y} - A_{9}e^{-m_{3}y} - A_{10}e^{-m_{1}y}\right)$$

$$u_{1}(y) = u_{10}(y) + Eu_{11}(y)$$

$$u_{1}(y) = \left(A_{17}e^{-m_{7}y} + A_{12}e^{-m_{6}y} - A_{13}e^{-m_{4}y} - A_{10}e^{-m_{1}y}\right)$$

$$+E\left(A_{25}e^{-m_{8}y} + A_{18}e^{-m_{7}y} + A_{19}e^{-m_{6}y} + A_{20}e^{-m_{5}y} - A_{24}e^{-m_{1}y}\right)$$

$$u_{1}(y) = \left(A_{25}e^{-m_{8}y} + A_{18}e^{-m_{7}y} + A_{19}e^{-m_{6}y} + A_{20}e^{-m_{5}y} - A_{24}e^{-m_{1}y}\right)$$

$$u_{1}(y) = \left(A_{21}e^{-m_{7}y} + A_{22}e^{-m_{6}y} - A_{23}e^{-m_{5}y} - A_{24}e^{-m_{1}y}\right)$$

$$u_{1}(y) = u_{0}(y) + \varepsilon e^{m_{1}y}(y)$$

$$C(y,t) = e^{-m_1 y} + \varepsilon e^{nt} \left(A_2 e^{-m_2 y} + A_1 e^{-m_1 y} \right)$$
(30)

$$\theta(y,t) = e^{-m_3 y} + \varepsilon e^{nt} \left(A_4 e^{-m_4 y} + A_3 e^{-m_3 y} \right)$$
(31)

Non-dimensional skin friction coefficient Cf, heat transfer rate and mass transfer rates are

$$\tau = \left(\frac{\partial u}{\partial y}\right)_{y=0} = \begin{bmatrix} \left(-m_5A_7 + m_3A_5 + m_1A_6\right) \\ +E\left(-m_6A_{11} - m_5A_8 + m_3A_9 + m_1A_{10}\right) \end{bmatrix}$$

$$-\varepsilon e^{nt} \begin{bmatrix} \left(-m_7A_{17} - m_5A_{12} + m_4A_{13} + m_3A_{14} \\ +m_2A_{15} + m_1A_{16} \\ +E\left(-m_8A_{25} - m_7A_{18} - m_6A_{19} - m_5A_{20} + m_4A_{21}\right) \end{bmatrix}$$
(32)

$$Nu = -\left(\frac{\partial\theta}{\partial y}\right) \text{at } (y=0) = (m_3) + \varepsilon e^{nt} (m_4 A_4 + m_3 A_3)$$
(33)

$$Sh = -\left(\frac{\partial C}{\partial y}\right) \text{at } (y=0) = (m_1) + \varepsilon e^{nt} (m_2 A_2 + m_1 A_1) \qquad (34)$$

4. RESULTS AND DISCUSSIONS

Selected computations have been depicted graphically in all the figures by assigning the values to the pertinent parameters characterizing the fluid flow mechanism. Extensive analytical computations are done for velocities, thermal and concentration distributions together with friction factor feature, Nusselt as well as Sherwood number for distinct standards of physical constraints which illustrate the structures of flow. Numerical conclusions are well established in Figures 2 to 10 additionally Tables 1 to 3.



Figure 2. Distribution of *u* for *Gr* & *Gm*

Figure 2, illustrates the impact of the flow for heterogeneous values of Grashof number Gr and modified Grashof number Gm. The figure conveys that, Grashof number Gr, modified Grashof number Gm are enlarged, it is perceived that the velocity elevated in general. Also, it is observed that as we shift away from the plate it is perceived that the effect of Grashof number Gr, modified Grashof number Gm are spotted to be not that notable. Figure 3, illustrates the end result of magnitic field parameter M and permeability parameter K on velocity profile. It is fascinating to note from figure that the

repercussion of magnetic field is to slow down the value of the velocity profile. The spire value radically declines with raise in the value of the magnetic field, because, the existence of magnetic field incites a force called the Lorentz force. It is also noticed from Figure 3 that the velocity of the flow accelerates with the increase of permeability parameter K.



Figure 3. Distribution of *u* for *M* & *K*



Figure 4. Distribution of *u* for *Pr*

Figures 4 and 8 display the dimensionless flow of the fluid and thermal boundary profiles for different values of Pr. These figures explore that the fluid flow and thermal boundary deescalate with uplift in Pr. This is due to the fact that with the raising values of Pr, thermal conductivity decreases, thus the velocity and temperature decreases with an increase in Pr.

Figures 5 and 9, show the antecedent profiles for assorted values of heat sink parameter Q and radiation R. The fluid flow and temperature profiles decrease with the increase of Q and the fluid flow and temperature profiles are increased with the increase of radiation parameter R. Figure 6 displays the dimensionless velocity for different values of Schmidt number Sc. We observe in this figure that the velocity profiles are reduced with increasing values of the Schmidt number Sc. The influence of visco-eleastic parameter E and suction parameter A on velocity profiles has been illustrated in Figure 7. It can identify that when E and A increase the velocity profile increases. From Figure 10, it can be seen that concentration distribution is detraining function of Sc. Further, it is seen that Sc does not contribute much to the concentration field as we

move far away from the boundary surface. Analogous effect is noted with chemical reaction parameter Kr on concentration profile.



Figure 5. Distribution of *u* for *Q* & *R*



Figure 6. Distribution of *u* for *Sc*



Figure 7. Distribution of *u* for *E* & *A*



Figure 8. Distribution of T for Pr



Figure 9. Distribution of *T* for *Q* & *R*



Figure 10. Distribution of C for Sc & Kr

The effects of varous parameter such as Grashaf value Gr, modified Grashaf value Gm, magnitic parameter M, porosity parameter K, Prandtl value Pr, radiation parameter R, viscoelastic parameter E, suction parameter A, heat sink parameter Q, Schmidt value Sc and chemical reaction parameter Kr on the friction factor (Cf), Nusselt value (Nu) and Sherwood value (Sh) are represented in Tables 1 to 3. From the Table 1, It is perceptible that as Grashaf value Gr or modified Grashaf value Gm or porosity parameter K or suction parameter A enhances, the friction factor uplifts, where as the friction factor downtrends as magnetic parameter M or visco-elastic parameter E increases from Table 1. It is concluded that as radiation parameter R or heat sink parameter Q or Prandtl

number Pr escalates, it is observed from Table 2 that the friction factor and Nusselt number escalates. From Table 3, it is found that the Sherwood value escalate when both the Schmidth value Sc and chemical reaction parameter Kr accelerate.

Table 1. Impact of different physical parameter on friction factor, Nusselt value and Sherwod value for Pr = 0.7, Q = 0.0, R = 0.2, Sc = 0.1, Kr = 2.0

Gr	Gm	М	K	Ε	Α	Cf	Nu	Sh
8.0	1.0	0.2	1.0	0.01	0.1	2.0115	5.7375	0.6432
9.0						2.2585	5.7375	0.6432
10.0						2.5055	5.7375	0.6432
	2.0					2.1408	5.7375	0.6432
	3.0					2.2702	5.7375	0.6432
		0.4				3.6009	5.7375	0.6432
		0.6				3.3221	5.7375	0.6432
			1.5			5.5941	5.7375	0.6432
			2.0			7.8830	5.7375	0.6432
-				0.02		6.0344	5.7375	0.6432
-				0.03		4.0230	5.7375	0.6432
					0.15	3.0172	5.7375	0.6432
-					0.17	3.4195	9.7538	0.6432

Table 2. Impact of different physical parameter on friction factor, Nusselt value and Sherwod value for Gr = 8.0, Gm = K = 1.0, M = 0.2, Ec = 0.01, Kr = 2.0, Sc = A = 0.1.

Pr	R	Q	Cf	Nu	Sh
0.7	0.2	0.5	2.0115	5.7375	0.6432
0.8			2.8576	1.3118	0.6432
1.0			8.7732	3.3588	0.6432
	0.8		3.8078	1.8848	0.6432
	0.9		4.4935	1.9581	0.6432
		1.0	3.3065	1.8110	0.6432
		1.5	9.9184	2.1758	0.6432

Table 3. Impact of different physical parameter on friction factor, Nusselt value and Sherwod value for Gr = 8.0, Gm = 1.0, K = 1.0, M = 0.2, R = 0.2, Pr = 0.7, Ec = 0.01, Kr = 2.0.

Sc	Kr	Cf	Nu	Sh
0.1	2.0	2.0115	5.7375	0.6432
0.3		2.0313	5.7375	1.2076
0.5		2.1053	5.7375	1.6467
	3.0	2.0124	5.7375	0.7667
	4.0	2.0146	5.7375	0.8714

5. CONCLUSIONS

In the present study a mathematical model has been developed to simulate 2D unsteady magneto hydrodynamic flow of an incompressible electrically conducting fluid over a permeable moving plate through porous medium under the importance of thermal radiation and chemical reaction. The governed mathematical statement is handled analytically by perturbation technique. The obtained results have led to the following conclusions:

- Fluid velocity is enhancing function of all parameters such as Grashaf number *Gr*, modified Grashaf number *Gm*, Permeability parameter *K*, Radiation parameter *R*, visco–elastic parameter *E*, Suction parameter *A*.
- Thermal boundary distribution falls down against *Pr* and *Q*, while radiation parameter *R* enhances it.
- Presence of chemical reaction enhances the rate of mass transfer which is a desired consequence of the flow of reacting species.
- Friction factor downtrends when magnetic parameter is enlarged. Nussselt number enhances for huge values of Pr. By increasing Schmidt number or chemical reaction parameter Sherwod number progress.

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NOMENCLATURE

A

suction parameter

B_0	magnetic field intensity, N. m ⁻¹ .A ⁻¹
C^{*}	fluid concentration
C_{∞}	free stream dimensional concentration
Α	visco-elastic parameter
D^{*}	Brownian diffusion coefficient, m ² . s ⁻¹
Gm	modified Grashof number
Gr	grashof number
$K_{\lambda w}$	absorption coefficient at wall
k^*	permeable parameter
K K_1	non-dimensional porous parameter chemical reaction coefficient
Kr	chemical reaction parameter
M	magnetic parameter
n^*	constant
Nu	local Nusselt number
P^{*}	pressure, Pa
Pr	Prandtl number
Q_0^*	dimensional heat sink
Q	heat sink parameter
R	radiation parameter
R^{*}	rate of chemical reactive factor
Sc	Schmidt number
Sh	Sherwood number
T^{*}	Temperature, K
T_{∞}	surface temperature, K
T_{∞}	free stream dimensional temperature K
U	stretching velocity, m. s ⁻¹
U_0	reference velocity, m. s ⁻¹
<i>u</i> *, <i>v</i> *	velocity components in x^*, y^* directions,
	m. s ⁻¹
$\hat{u_p}$	wall dimensional velocity, m. s ⁻¹
<i>x</i> *, <i>y</i> *	Cartesian coordinate in horizontal and vertical directions

Greek symbols

	1.:
v	kinematic viscosity, m ² . s ²
ε	small value
σ	electrical conductivity of the fluid, S. m ⁻¹
ρ	fluid density, kg. m ⁻³
v_0	constant suction velocity
$ au_w$	surface shear rate, Pa
β_T	coefficient of thermal expansion
β_C	coefficient of solutal expansion
k	thermal conductivity of the fluid, W. m^{-1} . k^{-1}
c_p	specific at constant pressure, J.K ⁻¹ .kg ⁻¹

APPENDIX

$$m_{1} = \frac{Sc + \sqrt{Sc^{2} + 4ScKr}}{2}, m_{2} = \frac{Sc + \sqrt{Sc^{2} + 4Scn_{5}}}{2}, m_{3} = \frac{\Pr + \sqrt{\Pr^{2} + 4\Pr n_{3}}}{2}, m_{4} = \frac{\Pr + \sqrt{\Pr^{2} + 4\Pr n_{4}}}{2},$$

$$\begin{split} m_5 &= \frac{1 + \sqrt{1 + 4n_1}}{2} , \quad m_6 = \frac{1 + \sqrt{1 + 4n_1}}{2} , \quad m_7 = \frac{1 + \sqrt{1 + 4n_2}}{2} , \qquad A_1 \\ m_8 &= \frac{1 + \sqrt{1 + 4n_2}}{2} , \quad A_1 = \frac{AScm_1}{m_1^2 - Scm_1 - Scn_5} , \quad A_2 = 1 - A_1 , \qquad A_1 \\ A_3 &= \frac{A\Pr m_3}{m_3^2 - \Pr m_3 - \Pr n_4} , \quad A_4 = 1 - A_3 , \quad A_5 = \frac{Gr}{m_3^2 - m_3 - n_1} , \qquad A_6 \\ A_6 &= \frac{Gm}{m_1^2 - m_1 - n_1} , \quad A_7 = u_p + A_5 + A_6 , \qquad A_8 \\ A_8 &= \frac{A_7 m_5^3}{m_5^2 - m_5 - n_1} , \quad A_9 = \frac{A_5 m_3^3}{m_3^2 - m_3 - n_1} , \quad A_{10} = \frac{A_6 m_1^3}{m_1^2 - m_1 - n_1} , \qquad A_{11} \\ A_{11} &= A_{10} + A_9 - A_8 \\ A_{12} &= \frac{AA_7 m_5}{m_5^2 - m_5 - n_2} , \quad A_{13} = \frac{GrA_4}{m_4^2 - m_4 - n_2} , \qquad A_{14} = \frac{GrA_3 + AA_5 m_3}{m_3^2 - m_3 - n_2} , \quad A_{15} = \frac{GrA_2}{m_2^2 - m_2 - n_2} , \qquad A_8 \end{split}$$

$$\begin{split} A_{16} &= \frac{GmA_1 + AA_6m_1}{m_1^2 - m_1 - n_2}, A_{17} = A_{16} + A_{15} + A_{14} + A_{13} - A_{12}, \\ A_{18} &= \frac{m_7^3 A_{17} + nA_{17}m_7^2}{m_7^2 - m_7 - n_2}, A_{19} = \frac{AA_{11}m_6}{m_6^2 - m_6 - n_2}, \\ A_{20} &= \frac{AA_7m_5^3 + AA_8m_5 + A_{12}m_5^3 + nA_{12}m_5^2}{m_5^2 - m_5 - n_2}, \\ A_{21} &= \frac{m_4^3 A_{13} + nA_{13}m_4^2}{m_4^2 - m_4 - n_2}, \\ A_{22} &= \frac{AA_5m_3^3 + AA_9m_3 + A_{14}m_3^3 + nA_{14}m_3^2}{m_3^2 - m_3 - n_2}, \\ A_{23} &= \frac{m_2^3 A_{15} + nA_{15}m_2^2}{m_2^2 - m_2 - n_2}, \\ A_{24} &= \frac{AA_6m_1^3 + AA_{10}m_1 + A_{16}m_1^3 + nA_{16}m_1^2}{m_1^2 - m_1 - n_2}, \\ A_{25} &= -A_{19} - A_{18} + A_{24} + A_{23} + A_{22} + A_{21} - A_{20}. \end{split}$$