



Heat Transfer Calculation on Viscous-Gravitational Fluid Flow Inside Vertical and Inclined Tubes

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ABSTRACT

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In this paper is presented a new procedure for the calculation of the heat transfer in vertical and inclines tubes, with presence of viscous-gravitational flow. Three models, one for inclines tubes and two for vertical tubes with presence of downward and upward flow were developed. The first model was development for coincident directions and 12 different fluids, included water and organic liquids, shows a correlation adjustment with a mean error of 12.75% in 81.51 % of the available experimental data in the interval of validity $7.5 \times 10^5 \leq Ra < 2.75 \times 10^{11}$ and $0.9 < Pr \leq 3.5 \times 10^4$. The second model was development for not coincident directions and 10 different fluids, included water and organic liquids, shows a correlation adjustment with a mean error of 13.04% in 83.09 % of the available experimental data in the interval of validity $7.6 \times 10^5 \leq Ra < 1.45 \times 10^{11}$ and $0.8 < Pr \leq 3.9 \times 10^4$. The third model was development for inclined tubes, it is valid for 10 different fluids included water and organic liquids, shows a correlation adjustment with a mean error of 16.12%, in 81.08% of the available experimental data in the interval of validity $7.9 \times 10^5 \leq Ra < 6.98 \times 10^{10}$, $0.9 < Pr \leq 2.1 \times 10^4$ and angle of inclination with respect to horizontal line $1^\circ \leq \theta \leq 88^\circ$. The objective of this paper is to make a procedure of analysis that enables considering the influence of the gravitational effects in the laminar flow regimen and to decrease the average uncertainty of the models and that additionally has a larger range of applicability.

1. INTRODUCTION

In the thermal engineering, the evaluation of the heat transfer processes inside of tubes, two models of not isothermic laminar flow are possible: viscous and viscous-gravitational. Each one of them has their own laws of heat transfer.

In laminar flow, when it comes true than $Ra \geq 8 \times 10^5$, then, the gravitational forces have influence of important way in the heat transfer, generating that the effect of the free convection cannot be rejected, therefore, the methods of analysis known for laminar flow does not allow obtaining satisfactory results. This type of problems is known as regimen of viscous-gravitational flow [1].

The viscous flow not isothermic appear when the forces of viscosity prevail over the gravitational forces, therefore, the viscous fluids do not receive the influence of the free convection. The viscous-gravitational flow appears when the gravitational forces are appreciable, and therefore, the natural convection adds to the forced. The model of viscous flow is so much more probable than the gravitational one, as much as the minor is the diameter of the tube, bigger the viscosity of the fluid and minor the temperature difference.

The velocity distribution in the cross section of the tube with viscous flow is not parabolic, due to the change of the viscosity with the temperature in the straightaway section. The velocity distribution is dependent if the fluid be heated or cooled (see Figure 1) [2].

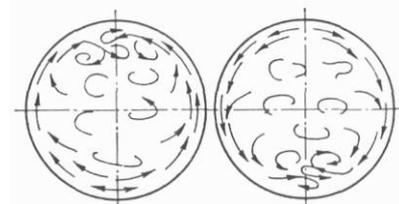


Figure 1. Transverse circulation in a tube due to the free convection (heating and cooling)

For a same mean temperature in the cross section, the temperature for the fluid near of the wall is major when he gets heat, that when he gets cold. While the temperature of a liquid increases, viscosity decreases, therefore, the velocity of the fluid near of the tube wall, in the case of heating, is greater than in that of cooling, and therefore, velocity increases in the heat transfer [3, 4].

When the liquid heats, the superficial transmission factor of heat is major that when he gets cold, the difference between the coefficients is major as much as the difference of temperature is bigger. When the liquid is heat, the superficial transmission factor of heat is major that when he gets cold; The difference between the heat transfer coefficients is major as much as bigger is the temperature difference [5].

In addition to the effect of the viscosity variations, in a gravitational viscous flow the distribution of velocity is very influenced for intensity and the direction of the free convection,

that originates as a consequence of the difference of densities between the hot and cold zones of the fluid [6, 7].

Therefore, the presence of the gravitational forces in the laminar flow regimen, makes complex the solution of the problem and reduces the precise grade in the obtained results, because the existing methods in literature do not consider this influence. Several investigators accomplished efforts to solve this inconvenience, however, the methods proposed only comprise reduced zones and close errors to the 30% [3].

Therefore, the authors follow as objective in this paper, developing a procedure of analysis that it enables considering the influence of the gravitational effects in the laminar flow regimen, that reduces the average level of uncertainty of the model and that additionally has a larger range of applicability.

2. METHODS AND VALIDATION

2.1 Introductory elements on the gravitational viscous flow

In the practical engineering, three basic cases of viscous gravitational regimen can be found, which depend on the direction of the forced and free convection [3]:

Case 1: The natural and the forced convection have the same direction.

Case 2: The natural and the forced convection have perpendicular directions.

Case 3: The natural and the forced convection have opposed directions.

Case 1 is found in vertical tubes, when a fluid that move along the tube in upward direction while receive heat, or a fluid that circulates in downward direction and its cooling. In this case, the effect of the free convection produces an increase of the fluid velocity near of the wall (see Figure 2) and there can be two maximums in the diagram of velocities distribution [8, 9].

In Figure 2 below, the enumerated curves are:

- 1) curve summation
- 2) forced convection
- 3) free convection

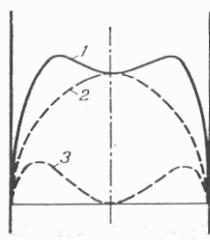


Figure 2. Distribution of velocities in the cross section of a tube with forced and free convection in the same direction

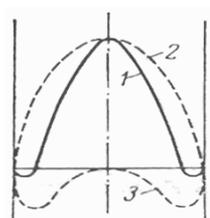


Figure 3. Distribution of velocities in the cross section of a tube with forced and free convection in the contrary direction

Case 2 is found in horizontal pipes. Due to the free convection, in the normal section of the tube a cross-sectional circulation comes from the fluid. In a heated fluid appear currents of free convection, upwards for the side of the wall and downwards for the center of the tube, while in the case of fluid cooling, the opposite happens. In consequence, the fluid moves through the tube following a spiral. The velocity of heat transfer increases due to the improvement in the fluid mixed [10, 11].

Case 3 is found in vertical tubes, when a fluid that move along the tube in downward direction while receive heat, or a fluid that circulates in upward direction and its cooling. In this case, decrease the velocity of the fluid near the wall, because convection currents have opposed directions (see Figure 3), which generates the formation of vortexes in the fluid close to the wall. Therefore, the process intensifies and the heat transfer turns out to be bigger than in the two previous cases, due to the intermittent appearing of turbulent motion [12, 13].

At the present time, in the specialized literature, the existing methods do not cover up with enough exactness, the estimation of the heat transfer in viscous-gravitational regimen in inclined tubes. A criterion that enjoys certain popularity is Cebeci's method, however, his use looks limited to inclined tubes 30° with respect to horizontal line, for what flows water in downward direction [4, 5].

In Figure 3 above, the enumerated curves are:

- 1) curve summation
- 2) forced convection
- 3) free convection

By means of the dimensional analysis techniques, it is very easy to prove that the viscous –gravitational flow regimen is strongly influenced by Prandtl and Grashoff dimensionless numbers, the inner diameter of the tube and its length [14].

The investigations realized in the years 50 of the last century focused on the solution of this problem, however, in spite of the great quantity of experimental data generated, the obtained solutions are based on systems of the mathematical form $Nu = aRe^m Pr^n$, for what it was required to limit the expressions obtained to intervals reduced of applicability [15].

However, the viscous–gravitational flow regimen in horizontal pipes is largely studied; existing in the literature an important group of current contributions, for such motive is excluded of the interests and reaches of this paper.

For the analysis of the cases 1 and 3, in literature exist two procedures, however these show high errors of correlation and have reduced ranges of applicability. These methods come given for [16-18]:

For the case 1

$$Nu = 0.35Gz^{0.3} \left[Ra \frac{d}{l} \right]^{0.18} \quad (1)$$

Eq. (1) is valid for the following range of values:

$$20 < l/d \leq 130 ; Re \leq 7.26Ra^{0.4}$$

$$8 \times 10^5 \leq Ra \leq 4 \times 10^8 ; 1.5 \left(Ra \frac{d}{l} \right)^{0.25} \leq Gz \leq 110$$

Eq. (1) correlates with a 50% of average error.

For the case 3

$$Nu = 0.037Re^{0.75} Pr^{0.4} (\mu_F / \mu_P)^N \quad (2)$$

In Eq. (2), the constant N take values 0.11 and 0.25 for heating and cooling of the fluid respectively.

Eq. (2) is valid for the following range of values:

$$0.2 < Pr \leq 100 ; 250 < Re \leq 10^4 \\ 1.5 \times 10^6 < Re \leq 12 \times 10^6$$

The Eq. (2) correlates with a 40% of average error.

3. DEVELOPMENT OF THE PROPOSAL MODELS

3.1 Deduction and validation of the proposal models

In the present paper, the experimental available data were obtained from an extended revision [3] compiled from literature specialized on the matter in Russian language [6-15], in which, were compiled and detailed an important group of experimental measurements on viscous gravitational flow, accomplished for investigators of the ancient Soviet Union. These experimental data will be used to achieve the adjustment and validation of the proposed models [19-21]. The study executed for the cases 1 and 3, as well as for inclined pipes, is given at once.

Case 1

Table 1 provides a summary of the experimental data used in the adjustment and validation of the proposal model for heat transfer calculation in viscous-gravitational flow for the case 1 [3]. The adjustment of the experimental available data is possible to realize it by means of a function of superposition (Brezhnetsov's function), establishing as fixed parameter the function $Gr^{0.2}$ and as residue $x^{0.48}$, in this case the residue is contingent upon the Prandtl number, in order to be the second experimental parameter in importance. The obtained model is given by:

$$Nu_{v1} = \frac{0.69Pr^{0.48}Gr^{0.2}(l/d)^{0.01}}{(0.61 + 1.2(Pr + 1.24Pr)^{0.48})^{0.3}} + A \quad (3)$$

In Eq. (3), if $Pr \leq 10$, then $A = 0$. For $Pr > 10$, then this should be corrected by means of the addition of a constant, whose value can be obtained in the Table 2.

Table 3 is given the validity intervals of the Equation (3). The Table 4 shows the correlation index of the Equation (3) in eight sub-intervals of their validity range, being proved that the Eq. (3) provides a correlation adjustment with an average error of 12.75% in 81.51% of the experimental available data. Figure 4 represents the adjustment obtained (with a 15% error band) between Eq. (3) and the experimental data.

Case 3

Table 5 provides a summary of experimental data used in the adjustment and validation of the proposal model for heat transfer calculation in viscous-gravitational flow for case 3.

The adjustment of the experimental available allow to get that the heat transfer in viscous-gravitational regimen in case 3 can be obtained as:

$$Nu_{v3} = \left[\frac{6.5GrPr^2}{100 + 105Pr} \right]^{0.2} + \frac{1090 + 1270Pr l}{2250 + 2200Pr d} + B \quad (4)$$

In Eq. (4), if $Pr \leq 10$, then $B = 0$. For $Pr > 10$, then this should be corrected by means of the addition of a constant, whose value can be obtained in the Table 6. In the Table 7 is given the range validity of Eq. (4). Table 8 shows the correlation index of Eq. (4) in eight sub-intervals of their validity range, being proved that Eq. (4) provides a correlation adjustment with an average error of 12.95% in 83.09% of the experimental available data. Figure 5 represents the adjustment obtained (with a 15% error band) between the proposed model and the experimental data.

Table 1. Experimental data used in the correlation of proposal model for the case 1

Source	Number of data	Fluid	$Ra \times 10^7$	Pr	l/d	Deviation percent
Petukhov (1950)	38	Water	0.075	0.8	10	16.3
			5.1	9.4	215	-10.5
Krasnochiekov (1957)	21	Water	0.21	1.2	20	4.4
			7.2	5.9	170	-5.1
Subbotin (1956)	17	Ethylene glycol	0.9	68	10	15.3
			1200	500	130	-11.8
Yashov (1960)	13	Dodecane	1.1	11	40	17.2
			150	28	180	-9.1
Yashov (1960)	9	Decane	1.4	7	30	11.2
			57	17	120	-4.2
Aladiev (1959)	31	MC Oil	1.2	130	40	17.2
			25000	9500	295	-14.1
Aladiev (1959)	26	MK Oil	0.092	580	32	16.3
			23420	35000	280	-13.7
Dodonov (1961)	23	Engine oil	1.2	90	40	18.1
			27520	21000	290	-14.7
Kern (1958)	5	Water	1.1	1.1	15	7.1
			6.2	5.2	125	-6.5
Boyko (1961)	41	Transformer oil	5.2	45.5	50	16.3
			4260	2950	120	-11.7
Ananiev (1962)	109	Water	0.078	2.1	20	6.1
			4.7	9.4	210	-7.2
Osipova (1963)	11	Methanol	0.095	2.3	30	5.2
			4.9	7.5	140	-11.8
Osipova (1963)	16	Ethanol	1.1	7.1	35	7.1
			3.2	60.2	150	-6.6

Mijeev (1953)	7	Glycerin	150	1830	50	19.3
			2400	18400	100	-18.4
Arafieva (1966)	17	Gasoline	3.2	5.8	30	5.1
			16.8	15.0	160	-13.8
For all sources above	384		0.075	0.8	10	13.1
			27520	35000	290	-11.9

Table 2. Values of the constant A used in Eq. (3)

Interval	Value of constant A
$7.5 \times 10^5 \leq Ra < 0.95 \times 10^8$	$0.27Ra^{0.25}$
$0.95 \times 10^8 \leq Ra < 1.1 \times 10^{10}$	$0.06Ra^{0.33}$
$1.1 \times 10^{10} \leq Ra \leq 2.75 \times 10^{11}$	$0.01Ra^{0.4}$

Table 3. Validity intervals to use the Eq. (3)

Parameter	Range
Fluids	Water, Ethylene glycol, Dodecane, Decane, MC oil, MK oil, Engine oil, Transformer oil, Methanol, Ethanol, Glycerin and Gasoline.
Pr	$0.8 \leq Pr \leq 3.5 \times 10^4$
Ra	$7.5 \times 10^5 \leq Ra \leq 2.75 \times 10^{11}$
l/d	$10 \leq l/d \leq 290$

Table 4. Correlation of Eq. (3) with experimental data

$7.5 \times 10^5 < Ra \leq 1 \times 10^6$	$0.9 < Pr \leq 10^2$	error < 8.69%
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$7.5 \times 10^5 < Ra \leq 5 \times 10^6$	$0.9 < Pr \leq 3 \times 10^2$	91.66% data error < 9.47%
$7.5 \times 10^5 < Ra \leq 1 \times 10^7$	$0.9 < Pr \leq 1.5 \times 10^3$	89.58% data error < 10.02%
$7.5 \times 10^5 < Ra \leq 5 \times 10^7$	$0.9 < Pr \leq 3.2 \times 10^3$	88.02% data error < 10.59%
$7.5 \times 10^5 < Ra \leq 1 \times 10^8$	$0.9 < Pr \leq 1.2 \times 10^4$	87.24% data error < 11.17%
$7.5 \times 10^5 < Ra \leq 1 \times 10^9$	$0.9 < Pr \leq 1.9 \times 10^4$	85.16% data error < 11.86%
$7.5 \times 10^5 < Ra \leq 1 \times 10^{10}$	$0.9 < Pr \leq 2.5 \times 10^4$	84.38% data error < 12.54%
$7.5 \times 10^5 < Ra \leq 2.8 \times 10^{11}$	$0.9 < Pr \leq 3.5 \times 10^4$	83.33% data error < 12.75%
		81.51% data

Table 5. Experimental data used in the correlation of proposal model for the case 3 [3]

Source	Number of data	Fluid	$Ra \times 10^7$	Pr	l/d	Deviation percent
Popov (1966)	43	Water	0.085	1.1	20	10.3
			4.3	7.2	160	-7.5
Nolde (1952)	36	Water	0.076	0.8	12	6.3
			11.2	9.9	250	-5.1
Liotsiansky (1955)	21	Benzene	0.9	3.1	10	7.2
			3.1	5	130	-4.2
Yashishtov (1962)	33	Kerosene	1.1	1.3	40	6.2
			15	2.9	180	-4.8
	11	Kerosene	1.4	1.4	20	3.7
			11.7	2.7	130	-5.2
Smirnova (1963)	17	Butanol	10.7	22.5	50	9.2
			1485	3810	270	-14.9
Petukhov (1957)	23	MK Oil	86.2	580	60	19.8
			13250	37100	260	-18.4
	19	MK Oil	113.1	650	60	21.7
			8748	17420	290	-23.8
	17	Engine oil	192.1	190	50	22.4
			14500	39200	220	-24.3
Maslov (1961)	22	Gasoline	3.1	5.8	40	12.3
			99.4	15.1	190	-14.2
Kurshatov (1955)	42	Water	0.092	0.9	15	7.1
			3.6	5.2	180	-3.2
Kirkmenko (1959)	26	Transformer oil	7.2	45.8	25	19.4
			1465	1280	160	-18.3
Klimenko (1969)	29	Water	1.1	1.0	30	4.8
			7.3	9.2	160	-8.2
Aladiev (1970)	33	Water	0.088	1.5	45	7.9
			6.2	8.3	190	-8.3
Godunov (1970)	11	Turpentine	7.8	14.1	30	14.2
			96.4	25.3	140	-16.2
Mijeev (1957)	31	Glycerin	125.1	2250	40	20.3
			3980	21940	180	-20.9
For all sources above	414		0.076	0.8	10	13.4
			14500	39200	290	-12.3

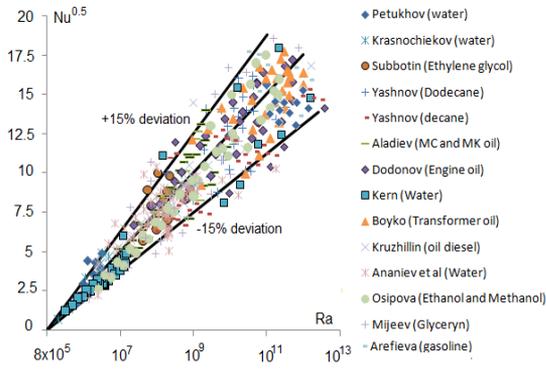


Figure 4. Adjust of the Eq. (3) with experimental data

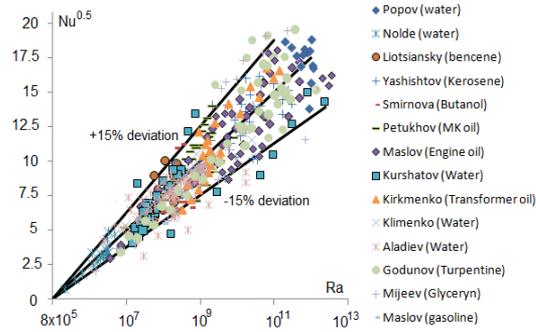


Figure 5. Adjust of the Eq. (4) with experimental data

Table 6. Values of the constant B used in Eq. (4)

Interval	Value of constant B
$7.6 \times 10^5 \leq Ra < 0.8 \times 10^9$	$0.38Ra^{0.26}$
$0.8 \times 10^9 \leq Ra \leq 1.45 \times 10^{11}$	$0.07Ra^{0.33}$

Table 7. Validity intervals to use the Eq. (4)

Parameter	Range
Fluids	Water, Benzene, Kerosene, Butanol, MC oil, MK oil, Engine oil, Transformer oil, Gasoline, Glycerin and Turpentine.
Pr	$0.8 \leq Pr \leq 3.92 \times 10^4$
Ra	$7.6 \times 10^5 \leq Ra \leq 1.45 \times 10^{11}$
l/d	$10 \leq l/d \leq 290$

Table 8. Correlation of Eq. (4) with experimental data

$7.6 \times 10^5 < Ra \leq 1 \times 10^6$	$0.8 < Pr \leq 1.2 \times 10^2$	error < 9.12% 90.88% data
$7.6 \times 10^5 < Ra \leq 5 \times 10^6$	$0.8 < Pr \leq 5 \times 10^2$	error < 9.81% 89.62% data
$7.6 \times 10^5 < Ra \leq 1 \times 10^7$	$0.8 < Pr \leq 1.3 \times 10^3$	error < 10.13% 88.41% data
$7.6 \times 10^5 < Ra \leq 5 \times 10^7$	$0.8 < Pr \leq 3.4 \times 10^3$	error < 10.61% 87.21% data
$7.6 \times 10^5 < Ra \leq 1 \times 10^8$	$0.8 < Pr \leq 0.8 \times 10^4$	error < 11.22% 85.74% data
$7.6 \times 10^5 < Ra \leq 1 \times 10^9$	$0.8 < Pr \leq 1.6 \times 10^4$	error < 11.77% 84.78% data
$7.6 \times 10^5 < Ra \leq 1 \times 10^{10}$	$0.8 < Pr \leq 2.7 \times 10^4$	error < 12.45% 83.57% data
$7.6 \times 10^5 < Ra \leq 1.5 \times 10^{11}$	$0.8 < Pr \leq 3.9 \times 10^4$	error < 13.04% 83.09% data

Table 9. Experimental data used in the correlation of proposal model for inclines tubes [3]

Source	Number of data	Fluid	$Ra \times 10^7$	Angle θ	Pr	l/d	Deviation percent
Ananiev (1956)	57	Water	0.082	30	0.9	15	14.3
			4.9	40	7.7	170	-13.5
Krasnochiekov (1961)	52	Water	0.079	65	1.2	20	13.2
			3.7	70	8.5	150	-11.8
Ribatsky (1956)	26	Kerosene	0.4	75	1.4	30	10.4
			9.1	88	2.8	230	-9.8
Yanishesky (1958)	25	Dodecane	0.9	15	11.2	50	16.4
			19.8	30	27.5	180	-14.8
Krasnov (1962)	19	Propanol	9.2	25	23	40	15.4
			22.4	40	30	150	-16.2
Eliazarov (1963)	67	Water	0.1	45	1.4	30	12.4
			6.9	60	9.2	140	-16.1
Malukshentov (1966)	47	Glycerin	29.4	30	1710	40	19.4
			7002	60	21002	190	-20.2
Udalov (1961)	32	Water	0.16	60	0.9	40	8.2
			8.6	75	9.6	260	-13.5
Klimenko (1973)	41	Water	0.24	1	1.6	50	6.2
			9.2	25	7.2	280	-4.6
Ivanisevich (1970)	37	Olive oil	91.2	60	650	60	17.6
			412.3	75	820	210	-15.8
Kurtakervich (1972)	18	Ciclohexane	1.3	75	11.2	30	13.9
			32.8	85	19.5	170	-16.4
Volkoba et al. (1967)	17	Aniline	1.5	40	11.9	40	19.4
			19.6	60	110	165	-22.3
Aladiev (1968)	19	Butyl Alcohol	16.2	30	23	30	21.7
			88.7	60	30	220	-20.9
Alexeev (1967)	24	Pentane	1.8	2	4.5	40	19.6
			7.2	30	7.1	250	-18.7
For all sources above	481		0.079	1	0.9	10	15.8
			6990	88	21050	290	-16.2

Inclined Tubes

Table 9 shows a summary of experimental data used in the adjust and validation of the obtained model for heat transfer calculation in viscous-gravitational flow for inclined tubes.

The adjustment of the experimental available data allows to get that the heat transfer in viscous-gravitational regimen for inclined tubes can be obtain as:

$$Nu_i = 0.12Ra^{0.3+0.01 \sin \theta} + \frac{1.03}{\left[\left(\frac{d}{l}\right) Ra^{0.2}\right]^{0.5}} + C \quad (5)$$

In the Eq. (5) if $Pr \leq 8$, then $C = 0$. For $Pr > 8$, then this should be corrected by means of the addition of a constant, whose value can be obtained in the Table 10. In the Table 10 the positive sign is taken if there exist coincidence between the gravitational forces and the flow direction, otherwise the negative sign is taken. In the Table 11 is given the validity intervals for eight sub-intervals of the Eq. (5).

Table 10. Values of the constant C used in Eq. (5)

Inclination of the tube	Value of constant C
$1^\circ \leq \theta < 12^\circ$	$\sqrt[3]{Nu_{V2}} \mp 0.12\sqrt{Nu_{V1}}$
$12^\circ \leq \theta < 30^\circ$	$0.1\sqrt{Nu_{V1}} \mp \sqrt[4]{Nu_{V2}}$
$30^\circ \leq \theta < 60^\circ$	$\sqrt[3]{(Nu_{V1} \mp Nu_{V2})^2}$
$60^\circ \leq \theta \leq 88^\circ$	$\sqrt[4]{0.16 Nu_{V1} \mp Nu_{V2} ^3}$

Table 11. Vality intervals to use the Eq. (5)

Parameter	Range
Fluids	Water, Kerosene, Dodecane, Propanol, Glycerin, Olive oil, Cyclohexane, Aniline, Butyl alcohol, Pentane.
Pr	$0.9 \leq Pr \leq 2.1 \times 10^4$
Ra	$7.9 \times 10^5 \leq Ra \leq 6.98 \times 10^{10}$
l/d	$10 \leq l/d \leq 290$
θ	$1^\circ \leq \theta \leq 88^\circ$

Table 12. Correlation of the Eq. (5) with experimental data

$7.9 \times 10^5 < Ra \leq 1 \times 10^6$	$0.9 < Pr \leq 1.2 \times 10^2$	error < 11.88%
		89.39% data
$7.9 \times 10^5 < Ra \leq 5 \times 10^6$	$0.9 < Pr \leq 5 \times 10^2$	error < 12.71%
		88.35% data
$7.9 \times 10^5 < Ra \leq 1 \times 10^7$	$0.9 < Pr \leq 1.3 \times 10^3$	error < 13.16%
		87.52% data
$7.9 \times 10^5 < Ra \leq 5 \times 10^7$	$0.9 < Pr \leq 3.4 \times 10^3$	error < 13.81%
		86.91% data
$7.9 \times 10^5 < Ra \leq 1 \times 10^8$	$0.9 < Pr \leq 7.5 \times 10^3$	error < 14.37%
		85.24% data
$7.9 \times 10^5 < Ra \leq 1 \times 10^9$	$0.9 < Pr \leq 1.1 \times 10^4$	error < 14.99%
		83.99% data
$7.9 \times 10^5 < Ra \leq 1 \times 10^{10}$	$0.9 < Pr \leq 1.7 \times 10^4$	error < 15.61%
		82.32% data
$7.9 \times 10^5 < Ra \leq 7 \times 10^{10}$	$0.9 < Pr \leq 2.1 \times 10^4$	error < 16.12%
		81.08% data

Table 12 shows the correlation index of the Eq. (5) in eight sub-intervals of their validity range, being proved that the Eq. (5) provides an correlation adjustment with an average error of 16.12% in 81.08% of the experimental available data, while, in the Figure 6 is represents the adjustment obtained (with a 15% error band) between the proposed model and the experimental data.

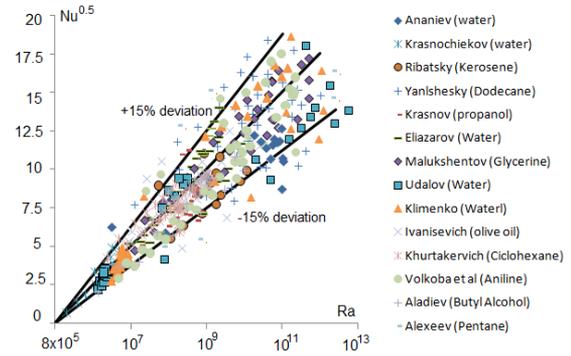


Figure 6. Adjust of the Eq. (5) with experimental data

4. CONCLUSIONS

Three new models have been development for heat transfer analysis in the viscous-gravitational flow. The proposal models show a bigger range of validity and a smaller error of correlation than the similarity models available in literature.

The first model is valid for the case 1 and is described by means of the Eq. (3). This model is valid for 12 different fluids, included water and organic liquids, shows a correlation adjustment with a mean error of 12.75% in 81.51% of the available experimental data in the interval of validity $7.5 \times 10^5 \leq Ra < 2.75 \times 10^{11}$ and $0.9 < Pr \leq 3.5 \times 10^4$.

The second model is valid for the case 3 and is described by means of the Eq. (4). This model is valid for 10 different fluids, included water and organic liquids, shows a correlation adjustment with a mean error of 13.04% in 83.09% of the available experimental data in the interval of validity $7.6 \times 10^5 \leq Ra < 1.45 \times 10^{11}$ and $0.8 < Pr \leq 3.9 \times 10^4$.

The third model is valid for inclined tubes and is described by means of the Eq. (5). This model is valid for 10 different fluids, included water and organic liquids, shows a correlation adjustment with a mean error of 16.12%, in 81.08% of the available experimental data in the interval of validity $7.9 \times 10^5 \leq Ra < 6.98 \times 10^{10}$, $0.9 < Pr \leq 2.1 \times 10^4$ and angle of inclination with respect to horizontal line $1^\circ \leq \theta \leq 88^\circ$.

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NOMENCLATURE

A	Constant, defined in Eq. (3).
B	Constant, defined in Eq. (4).
C	Constant, defined in Eq. (5).
d	Equivalent inner tube diameter, m
Gr	Grashoff number
Gz	Graetz number
l	Length of the tube, m
N	Constant, defined in Eq. (4)
Nu_{V1}	Nusselt number for vertical tubes, case 1
Nu_{V3}	Nusselt number for vertical tubes, case 3
Nu_i	Nusselt number for inclined tubes
Pr	Prandtl number
Ra	Rayleigh number
Re	Reynolds number
T_F	Average fluid temperature, °C
T_P	Wall temperature, °C

Greek symbols

α	Heat transfer coefficient, $\text{kg}\cdot\text{m}^{-1}\cdot\text{K}^{-1}\cdot\text{s}^{-1}$
μ_F	Fluid dynamic viscosity at T_F , $\text{kg}\cdot\text{m}^{-1}\cdot\text{s}^{-1}$
μ_P	Fluid dynamic viscosity at T_P , $\text{kg}\cdot\text{m}^{-1}\cdot\text{s}^{-1}$
θ	Tube inclination respect to horizontal line