

Soret effect on transient magnetohydrodynamic nanofluid flow past a vertical plate through a porous medium with second order chemical reaction and thermal radiation

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ABSTRACT

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In the present article, an investigation on the impact of Soret effect on unsteady magnetohydrodynamic natural convection flow of a viscous, electrically conducting and incompressible nanofluid over a vertical plate through a medium filled with porous materials with the consideration of a second order chemical reaction, has been investigated. Three kinds of water-based nanofluids, comprising of aluminum oxide (Al_2O_3), titanium oxide (TiO_2) and silver (Ag) as nanoparticles, are chosen for this analysis. Governing equations accompanied by the boundary and initial conditions are changed into non-dimensional form and then they are solved by Crank-Nicolson type finite difference scheme. The impact of relevant flow parameters on nanofluid velocity, nanofluid temperature and species concentration are depicted graphically with comprehensive discussions whereas numerical findings for coefficient of skin friction, wall temperature gradient i.e. Nusselt number and wall concentration gradient are depicted in tabular form. It has been observed that Soret effect has the ability to enhance the species concentration. The investigation we have performed here, has various scientific and industrial applications.

1. INTRODUCTION

The concept of nanofluids was proposed by Choi [1] at Argonne National Laboratory (ANL). Nanofluids are proved to be the promising fluids for the improvement of heat transfer because of their extraordinary thermal conductivity property, where nano-sized particles (minor than 100 nanometers) are suspended uniformly in a base fluid. The type of nanoparticle used is a kind of big deal as far as the enhancement of thermal conductivity of the base fluid is concerned. In the recent past, nanofluids are used extensively in numerous scientific processes such as transportation, power generation, treatment of cancer-infected tissues, electronic and space cooling, etc. Different nanosized particles of metals and their oxides such as Cu , Ag , Al_2O_3 and CuO , carbides (SiC) or non-metals (graphite, carbon nanotubes) are frequently used as nanoparticles while preparing nanofluid. Selvakumar and Dhinakaran [2] studied the Krieger–Dougherty (K-D) model to describe the effective nanofluid viscosity based on PSD (particle size distribution) study. They analysed the effect of aggregation together with the interfacial layer formation in the K–D model in order to make it a possible nanofluids. Heris et al. [3] investigated the laminar flow of Al_2O_3 -nanofluid in a circular tube experimentally and showed that the coefficient of convective heat transfer of pure water was improved by 41% with only 3% volume fraction of Al_2O_3 nanoparticle. Fotukian and Esfahany [4] considered the influence of CuO -nanofluid in a round tube in a turbulent flow situation and observed that the coefficient of convective heat transfer was boosted by 25% when the concentration of cupric oxide is taken low. In the recent few years, the impact of magnetohydrodynamic (MHD) nanofluid flows together with mass and heat transfer have

attracted several researchers, engineers and scientists due to their numerous real world applications in contemporary technology and many industries such as in cooling of fission reactors, power generation in space craft, stellar and solar structures, medicine delivery processing, hypersonic wind tunnel experiments, liquid metals, MHD pumps and radio propagation. Inspired from such real-world applications, many researcher and engineers have carried out the research work on MHD, experimentally as well as theoretically. Turkyilmazoglu [5] analysed the flow of four different kinds of nanofluids viz. Cu -nanofluid, Ag -nanofluid, TiO_2 -nanofluid and CuO -nanofluid and found the analytical solution for magnetohydrodynamic slip and mass transport in nanofluids. Sheikholeslami et al. [6] analyzed the flow of CuO -nanofluid under the effect of magnetic field and found the enhancements in heat transfer for higher values of magnetic parameter. Sandeep and Reddy [7] discussed MHD Cu -nanofluid flow near a cone and a wedge and they observed that volume fraction of nanoparticle and magnetic parameter strengthen the temperature field.

While analyzing the free convection flow of hydromagnetic liquids to characterize the heat transfer behavior, the radiative property of the liquid cannot be overlooked as it has substantial influence on the heat transfer characteristic of flow. This radiative property of the liquid affects the heat transfer performance of several industrial and engineering processes such as furnace design, turbine blade heat transfer, polymer processing, solar energy collector, glass production, heating and cooling chambers etc. Considering all these applications of thermal radiation effect in real flow situations, several scientists and researchers continue to delve deeper into the study of the influence of thermal radiation. Takhar et al. [8]

scrutinized the radiation effect on free convection flow of a hydromagnetic fluid over a semi-infinite vertical plate. Das and Jana [9] examined the MHD and thermal radiation effects on free convection flow of water based nanofluid with three kinds of nanoparticles, past an infinite moving horizontal plate. They deduced that a surge in the radiation corresponds to a deceleration in the fluid velocity. Some other important studies on the topic are due to Ogulu and Makinde [10], Mahmoud [11], Das et al. [12], Seth et al. [13-14].

It is observed that in all above experimental or theoretical studies, thermophoresis or Soret effect was ignored in mass transfer process because of their small order of magnitude. Thermophoresis or Soret effects are noteworthy when there exists a difference in the density in the flow field, as in the case of geo-sciences and hydrological fluid flow situation. Soret effect is observed in various flow situations such as, in a compound with medium molecular weight (N_2) and of very light molecular weight (H_2), in isotope separation, etc. For different fluids, many scientists have studied the thermophoresis effect due to its significant applications in science. Considering the significance of such effect, Ahmed [15] analysed the Hartmann flow restricted by two parallel isothermal horizontal plates and presented the analytical solution with the consideration of Soret effect. Hussanan et al. [16] examined the MHD mixed-convection flow past an oscillating vertical plate, in consideration of thermophoresis effect and observed that the species concentration gets enhanced with an increment in Soret number. Seth et al. [17] discussed the Soret effect on transient MHD natural convection flow near a vertical moving plate with rampedness in the temperature of the plate, in a rotating system. They observed that thermophoresis effect has a tendency to distribute the species throughout the flow boundary layer and it decreases the concentration gradient at the plate. Zhao et al. [18] considered the free convection stream of hydromagnetic fractional viscoelastic fluid in a porous medium, taking Dufour and Soret effects into account. Some other related research work on this theme are also due to Reddy and Chamkha. [19], Saritha et al. [20], Bég et al. [21], Abdel-wahed and S. M. Abdel-AAL [22].

Chemical reaction effect on mass transfer flows are vital in a lot of industrial applications. Some of the industrial and technological applications of chemical reactions are observed in chemical industries viz. polymer production, manufacture of ceramics, food processing, catalytic chemical reactors, hydrometallurgical industries, polymer production etc. Due to the significance of reaction of chemicals in all sorts of chemical processes in industries, several scientists have studied transient hydromagnetic free convection flow with higher order chemical reaction under different conditions. Das et al. [23] analysed the impact of chemical reaction on the fluid flow with uniform heat flux, near an infinite vertical plate and noticed that species concentration reduces with a surge in the rate of 1st order chemical reaction. Seth et al. [24] examined the free convection MHD flow past an impulsively moving plate with induced magnetic field and nth order chemical reactions where the governing equations were solved numerically by a finite difference scheme. They observed that chemical reaction's order has a tendency of enhancing the species concentration. Venkateswarlu et al. [25] inspected the influence of chemical reaction on hydromagnetic natural convection Jeffrey fluid in the vicinity of a permeable moving vertical plate with the consideration of heat source.

The present study seeks to explore the effect of thermo-

diffusion phenomena on the transient hydromagnetic free convection flow of an incompressible, viscous, optically thick heat radiating, chemically reactive and electrically conducting fluid near a vertical infinite plate through a porous medium with second order homogenous chemical reaction. Three varieties of nanoparticles, namely, aluminum oxide (Al_2O_3), silver (Ag) and titanium oxide (TiO_2) are chosen for this study. To the best of our information, no attempt has been made to undertake this kind of study. Such types of flow problem may arise in several industrial processes such as coolant application, manufacture of ceramics and glassware, nanodrug delivery, cooling of microchip, polymer production, etc.

2. MATHEMATICAL FORMULATION

Consider the transient magnetohydrodynamic free convection flow with the second order chemical reaction and Soret effect, of chemically reactive, optically thick heat radiating, incompressible, viscous and electrically conducting nanofluid near a vertical plate which is of infinite length and whose movement is impulsive, through a porous medium. Cartesian coordinate system is chosen to represent the geometry of the present model, such that the plate is considered to be in the positive direction of x' -axis. y' -axis is in a direction which is perpendicular to the plate whereas z' -axis is in the normal direction to $x'y'$ -plane. There is an application of an unvarying magnetic field of strength B_0 in a direction which is transversal to the motion of the plate. For $t' \leq 0$, i.e. in the initial phase, the plate as well as the fluid, both are kept at rest and are kept at constant temperature T'_∞ whereas level of concentration of nanoparticles is maintained at a constant value C'_∞ . The plate start moving impulsively in its own plane with velocity λU_0 in the positive direction of x' -axis at $t' > 0$, where U_0 is constant. The geometry of the present model is portrayed in Fig.1. Due to the small magnetic Reynolds number of metallic liquids, the influence of induced magnetic field has been ignored.

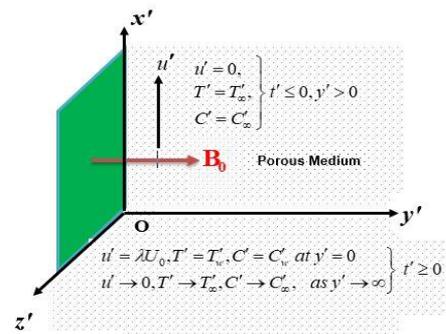


Figure 1. Geometry of the problem

Following the above made assumptions, the boundary layer equations governing the nanofluid velocity, temperature and distribution of nanoparticles, taking second order chemical reaction and Soret effect into account, under the Boussinesq approximation is given as follows:

$$\rho_{nf} \frac{\partial u'}{\partial t'} = \mu_{nf} \frac{\partial^2 u'}{\partial y'^2} - \left(\sigma_{nf} B_0^2 + \frac{\mu_{nf}}{k'} \right) u' + g (\rho \beta_C)_{nf} (C' - C'_\infty) + g (\rho \beta_T)_{nf} (T' - T'_\infty) \quad (1)$$

$$\frac{\partial T'}{\partial t'} = \frac{k_{nf}}{(\rho c_p)_{nf}} \frac{\partial^2 T'}{\partial y'^2} - \frac{1}{(\rho c_p)_{nf}} \frac{\partial^2 q'_r}{\partial y'^2} \quad (2)$$

$$\frac{\partial C'}{\partial t'} = D_{nf} \frac{\partial^2 C'}{\partial y'^2} + D_m \frac{\partial^2 T'}{\partial y'^2} - k'_2 (C' - C'_\infty)^2 \quad (3)$$

where $u', T', C', \rho_{nf}, k_{nf}, \sigma_{nf}, (\beta_T)_{nf}, (\beta_c)_{nf}, \mu_{nf}, q'_r, B_0, k', D_M$ and g , are fluid velocity, fluid temperature, species concentration, nanofluid density, thermal conductivity of nanofluid, electrical conductivity of nanofluid, thermal expansion coefficient of nanofluid, coefficient of expansion for species concentration, dynamic viscosity of nanofluid, radiative flux vector, uniform magnetic field, porous medium's permeability, thermal diffusion coefficient and acceleration due to gravity, respectively.

For nanofluid, the expressions of

$\rho_{nf}, (\rho\beta_T)_{nf}, (\rho\beta_c)_{nf}, \mu_{nf}, \sigma_{nf}, (\rho c_p)_{nf}$ and k_{nf} are given as:

$$\rho_{nf} = (1-\psi)\rho_f + \psi\rho_s, (\rho\beta_T)_{nf} = \psi(\rho\beta_T)_s + (1-\psi)(\rho\beta_T)_f,$$

$$(\rho\beta_c)_{nf} = \psi(\rho\beta_c)_s + (1-\psi)(\rho\beta_c)_f, \mu_{nf} = \frac{\mu_f}{(1-\psi)^{2.5}},$$

$$\sigma_{nf} = \sigma_f \left(1 + \frac{3(\sigma-1)\psi}{(\sigma+2) - (\sigma-1)\psi} \right),$$

$$(\rho c_p)_{nf} = \psi(\rho c_p)_s + (1-\psi)(\rho c_p)_f,$$

$$\frac{k_{nf}}{k_f} = \frac{-2\psi(k_f - k_s) + (k_s + 2k_f)}{\psi(k_f - k_s) + (k_s + 2k_f)},$$

where ρ_s is density of nanoparticles, ρ_f is the density of base fluid, μ_f is base fluid's viscosity, ψ is volume fraction of nanoparticle, $(\rho c_p)_f$ is the heat capacitance of base fluid, $(\rho c_p)_s$ is the heat capacitance of nanoparticles, σ_s is nanoparticles' electrical conductivity, k_s is nanoparticles' thermal conductivity and k_f is base fluid's thermal conductivity.

Following are the appropriate boundary and initial conditions:

$$u' = 0, C' = C'_\infty, T' = T'_\infty, \text{ for } y' > 0, \text{ and } t' \leq 0, \quad (4a)$$

$$\left. \begin{aligned} u' &= \lambda U_0, T' = T'_w, C' = C'_w \quad \text{at } y' = 0 \\ u' &\rightarrow 0, T' \rightarrow T'_\infty, C' \rightarrow C'_\infty \quad \text{as } y' \rightarrow \infty \end{aligned} \right\} t' \geq 0 \quad (4b)$$

Table 1. Thermo-physical properties of nanoparticles and water

	c_p ($Jkg^{-1}K^{-1}$)	k ($Wm^{-1}K^{-1}$)	σ (S/m)	$\beta \times 10^{-5}$ (K^{-1})	ρ (kgm^{-3})
Al_2O_3	765	40	1.07×10^6	0.85	3970
MoS_2	397.21	904.4	2.09×10^4	2.8424	5.06×10^{-3}
Ag	235	429	6.30×10^7	1.89	10500
H_2O	4179	0.613	5.5×10^{-6}	21	997.1

When the fluid is an optically thick radiating one, we can approximate the radiative flux q'_r by the expression proposed by the Rosseland. Therefore, according to the Rosseland approximation [26], for a medium which is considered optically thick, the radiative heat flux q'_r is given as:

$$q'_r = -\frac{4}{3} \left(\frac{\sigma^*}{k^*} \right) \frac{\partial T'^4}{\partial y'}, \quad (5a)$$

where k^* is Rosseland mean absorption co-efficient and σ^* is Stefan-Boltzmann constant. It is supposed that temperature difference inside boundary layer region is small enough such that the term T'^4 can be presented as a linear function of temperature T' . Therefore T'^4 may be expanded using Taylor's expansion about T'_∞ and by neglecting terms of order more than and equal to two in the expansion, we get

$$T'^4 \cong 4T'^3_\infty T' - 3T'^4_\infty. \quad (5b)$$

Making use of (5a) and (5b), equation (2) takes the following form:

$$\frac{\partial T'}{\partial t'} = \frac{k_{nf}}{(\rho c_p)_{nf}} \left(1 + \frac{16\sigma^* T'^3_\infty}{3kk^*} \right) \frac{\partial^2 T'}{\partial y'^2}. \quad (6)$$

To present equations (1), (3) and (6) together with boundary and initial conditions (4a) to (4b) in dimensionless form, we use following variables which are non-dimensional in nature:

$$\left. \begin{aligned} y &= y' / U_0 t_0, T = (T' - T'_\infty) / (T'_w - T'_\infty), u = u' / U_0, \\ C &= (C' - C'_\infty) / (C'_w - C'_\infty), t = t' / t_0 \end{aligned} \right\} \quad (7)$$

Using non-dimensional variables given by equation (7), equations (1), (3) and (6) get converted into following dimensionless form:

$$\frac{\partial u}{\partial t} = b_1 \frac{\partial^2 u}{\partial y^2} + b_2 Gr T + b_3 Gc C - Mb_4 u - b_4 \frac{u}{K}, \quad (8a)$$

$$\frac{\partial T}{\partial t} = b_5 \frac{\partial^2 T}{\partial y^2} \quad (8b)$$

$$\frac{\partial C}{\partial t} = \frac{1}{Sc} \frac{\partial^2 C}{\partial y^2} - Kr C^2 + Sr \frac{\partial^2 T}{\partial y^2}, \quad (8c)$$

The dimensionless boundary and initial conditions are presented as:

$$u = 0, T = 0, C = 0 \text{ for } t \leq 0 \text{ and } y \geq 0 \quad (9a)$$

$$\left. \begin{aligned} u &= \lambda, T = 1, C = 1 \quad \text{at } y = 0 \\ u &\rightarrow 0, C \rightarrow 0, T \rightarrow 0 \quad \text{as } y \rightarrow \infty \end{aligned} \right\} t \geq 0 \quad (9b)$$

where $Gr = \frac{\nu_f g \beta_T (T'_w - T'_\infty)}{U_0^3}$ is thermal Grashof number, $Nr = \frac{16\sigma^* T'^3_\infty}{3k_{1nf} k'}$ is radiation parameter, $Sc = \frac{\nu_f}{D_{nf}}$ is Schmidt number,

$Gm = \frac{\nu_f g \beta_{cf} (C'_w - C'_\infty)}{U_0^3}$ is solutal Grashof number, $Pr = \frac{(\mu c_p)_f}{k_{1f}}$ is Prandtl number, $Kr = \frac{\nu_f k_2 (C'_w - C'_\infty)}{U_0^2}$ is chemical reaction parameter, $M = \frac{\nu_f \sigma_f B_0^2}{\rho_{nf} U_0^2}$ is magnetic parameter and $Sr = \frac{D_M (T'_w - T'_\infty)}{\nu_f (C'_w - C'_\infty)}$ is Soret number.

The other parameters and variables appearing in equations (8a) to (8c) are given as:

$$b_1 = \frac{\mu_f}{(1-\psi)^{2.5}}, b_2 = \frac{x_2}{x_3}, b_3 = \frac{x_3}{x_1}, b_4 = \frac{x_6}{x_1},$$

$$b_5 = \frac{1}{x_4 Pr} (x_5 + Nr),$$

$$\text{where } x_1 = (1-\psi) + \psi \left(\frac{\rho_s}{\rho_f} \right), x_2 = \left[(1-\psi) + \psi \left(\frac{\rho \beta_T}{\rho \beta_T} \right)_s \right],$$

$$x_3 = \left[(1-\psi) + \psi \left(\frac{\rho \beta_c}{\rho \beta_c} \right)_s \right], x_4 = \left[(1-\psi) + \psi \left(\frac{\rho c_p}{\rho c_p} \right)_s \right],$$

$$x_5 = \frac{(k_s + 2k_f) - 2\psi(k_f - k_s)}{(k_s + 2k_f) + \psi(k_f - k_s)}, \frac{1}{K} = \frac{\nu_{nf} \psi \nu}{k U_0^2},$$

$$\text{and } x_6 = \left(1 + \frac{3(\sigma-1)\psi}{(\sigma+2) - (\sigma-1)\psi} \right).$$

3. NUMERICAL SOLUTION

Equations (8a) to (8c) are coupled, partial differential equations where equation (8c) is non-linear in nature, which therefore, cannot be solved analytically. Therefore, we solve these equations by implicit finite difference method of Crank-Nicolson type. This scheme corresponding to equations (8a) to (8c) with initial and boundary conditions (9a) and (9b) are applied in following manner:

The region is discretized as rectangular grid with grid lines parallel to the coordinate axes with placing Δy and Δt in space and time respectively. Also $\Delta y = y_{i+1} - y_i$ and $\Delta t = t_{j+1} - t_j$. With the known values of C, T and u at a time $t = j$, the finite difference equation corresponding to (8b) at each interval nodal point on a specific j -level creates a system of tridiagonal equations. This system is solved by the use of Thomas algorithm, as proposed by Carnahan et al. [27]. In this way, the T values are determined at every nodal point at $(j+1)^{\text{th}}$ time step. In the same manner, the C values are determined from equation (8c) using the T values at $(j+1)^{\text{th}}$ time step of equation (8c). Finally, with the use of C and T values at $(j+1)^{\text{th}}$ time step in equation (8a), the value of u at $(j+1)^{\text{th}}$ time step are obtained in similar fashion. This procedure is continued to obtain the solution until desired time t . Therefore, the values of T, C and u are determined at all points of the rectangular expanse for the desired time level. There is a truncation error of order $O\{(\Delta y)^2 + (\Delta t)^2\}$ in the Crank-Nicolson scheme which diminishes as Δy and Δt approach to zero, which justifies the consistency (Antia [28]). Consistency and stability together guarantee the convergence of scheme.

4. WALL VELOCITY GRADIENT, WALL TEMPERATURE GRADIENT AND CONCENTRATION GRADIENT IN THE VICINITY OF THE PLATE

Expression for wall velocity gradient τ , wall temperature gradient Nu and concentration gradient in the vicinity of the plate, S_h are calculated from equations (10) to (12) respectively and are presented as follows:

$$\tau = \frac{\partial u}{\partial y} \Big|_{y=0} \quad (10)$$

$$Nu = \frac{\partial T}{\partial y} \Big|_{y=0} \quad (11)$$

$$S_h = \frac{\partial C}{\partial y} \Big|_{y=0} \quad (12)$$

5. RESULTS AND DISCUSSION

To highlight the impact of several physical parameters which characterize the features of the flow, extensive numerical calculations have been performed for nanofluid velocity, temperature and species concentration together with wall velocity gradient, wall temperature gradient and concentration gradient near the plate, to get an insight of the physics involved in the flow regime.

Numerical results which are demonstrated in Figures 2-11, have been scrutinized in detail by giving numerical values to physical parameters. Throughout the study, we take the following values of non-dimensional numbers: $Gr = 5, Gc = 4, Nr = 4, M = 3, Sr = 2, Kr = 3, \varphi = 0.15, \gamma = 1, Pr = 0.71, Sc = 0.22, \lambda = 1$ and $K = 2$, unless stated otherwise.

The variation in the nanofluid velocity versus the boundary layer coordinate y , under the influence of various parameters are depicted graphically in Figures 2 to 6. Figures 2 and 3 portray the influence of Gr and Gc on nanofluid velocity u , respectively. It is conceived from these Figures that nanofluid velocity u increases on increasing the values of Gr and Gc . Since, an increase in Gr and Gc means an increase in the strengths of thermal and solutal buoyancy forces, respectively. This infers that solutal as well as thermal buoyancy forces tend to speed-up the fluid velocity.

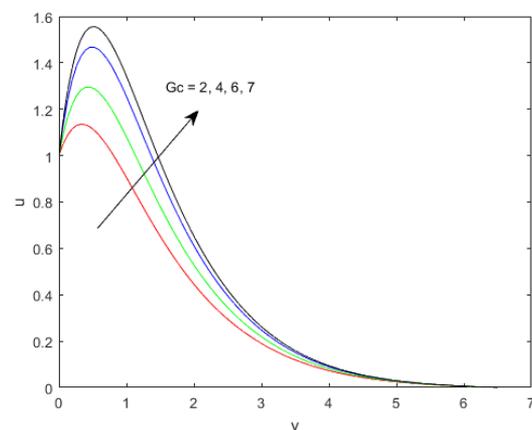


Figure 2. Velocity profiles for Gc

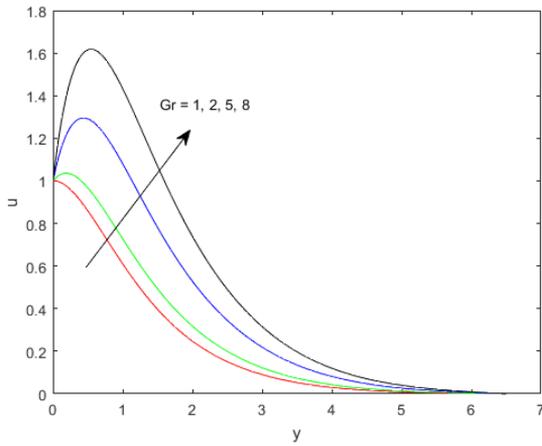


Figure 3. Velocity profiles for Gr

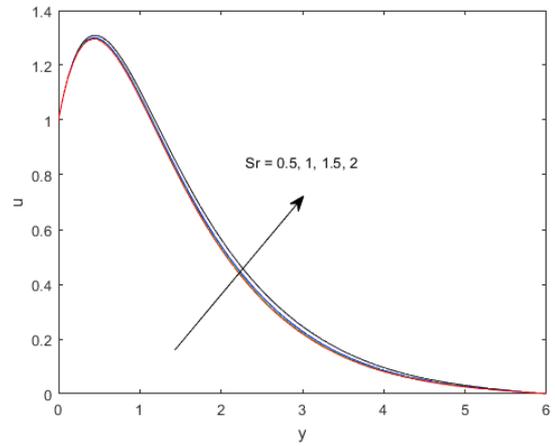


Figure 6. Velocity profiles for Sr

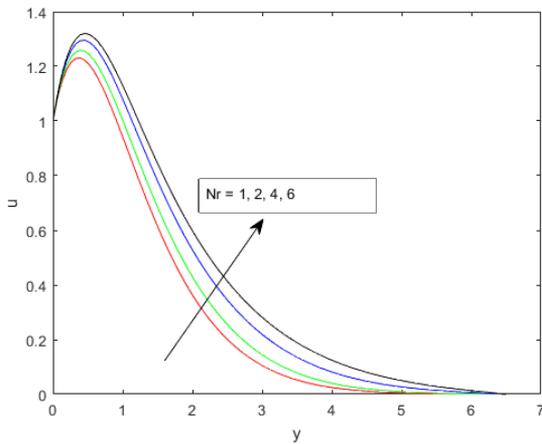


Figure 4. Velocity profiles for Nr

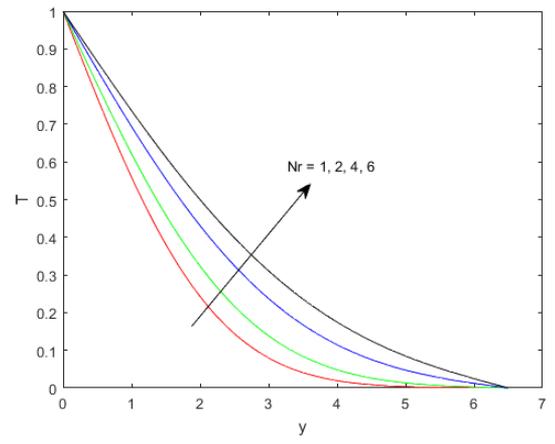


Figure 7. Temperature profiles for Nr

It is perceived from Figure 4 that, nanofluid velocity u is getting accelerated with the growing value of radiation parameter Nr . The reason behind such behavior of nanofluid velocity is that an intensification in radiation parameter outcomes in the increase of fluid temperature, and therefore thermal buoyancy force becomes stronger with amplified thermal radiation. This enhanced thermal buoyancy force results in quicker fluid motion.

It is noticed from Figure 5 that u is getting slowed down on strengthening the magnetic field. This fact is justified since the presence of magnetic field produces an opposing force, termed as Lorentz force, which works in a direction, reverse to the fluid course and has the tendency to resist the flow and hence a decelerated fluid motion is observed in this case.

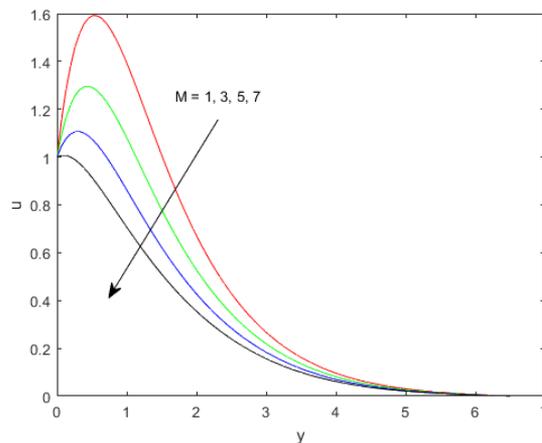


Figure 5. Velocity profiles for M

It is perceived from Figure 6 that nanofluid velocity u is increased with the growing value of Soret number Sr . This is in agreement with the fact that a rise in the value of Soret number enhances the solutal buoyancy force which results in accelerated fluid motion.

Figure 7 depicts, how the fluid temperature $T(y, t)$ is affected by the radiation parameter Nr . One can observe from this figure that temperature of nanofluid rises with the growing value of radiation parameter Nr .

The influence of nanoparticle volume fraction ϕ on temperature is depicted in the Figure 8, in which temperature is enhanced as the nanoparticle volume fraction is increased.

To demonstrate the characteristics of species concentration against chemical reaction, thermal-diffusion and mass diffusion parameters, the numerical solutions of species concentration $C(y, t)$ are obtained and are revealed in form of graphs in Figures 9 to 11. Figure 9 demonstrates concentration profile against Schmidt number Sc . It is noticed that species concentration C decreases on increasing the value of Sc . This phenomenon is in good agreement with the fact that a rise in Sc leads to a decrement in the molecular diffusion.

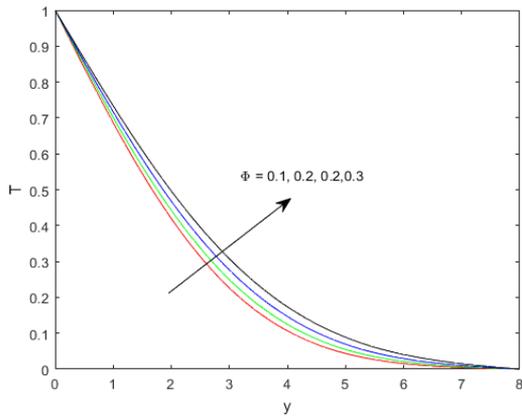


Figure 8. Temperature profiles for ϕ

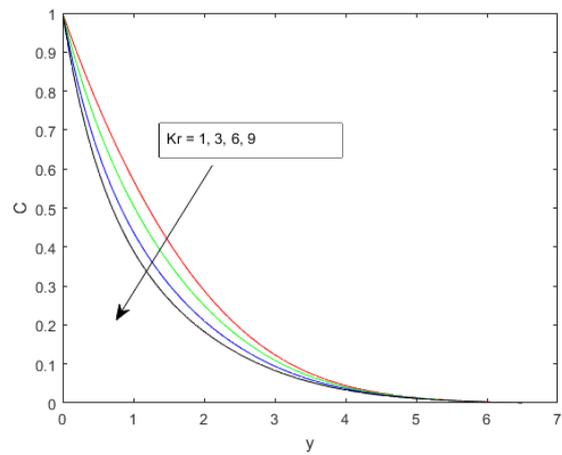


Figure 11. Concentration profiles for Kr

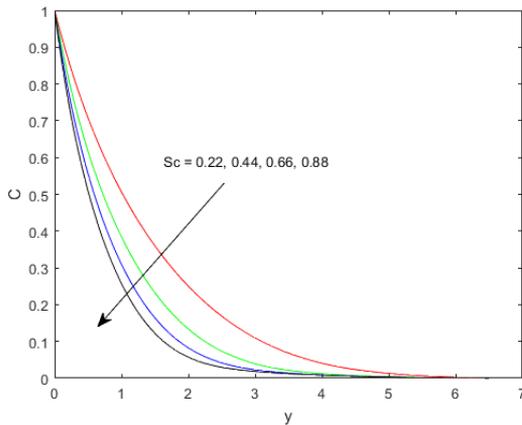


Figure 9. Concentration profiles for Sc

Highlights of the influence of Soret number Sr on concentration profile is represented in Figure 10. Concentration of species grows on increasing the value of Sr . This infers that Soret number tends to enhance species concentration.

It is evident from Figure 11 that concentration of species diminishes with the growing value of chemical reaction parameter Kr . The central reason behind this fact is that destructive kind of chemical reaction leads to a decline in the species concentration by virtue of which, the buoyancy effect due to concentration gradient gets nullified.

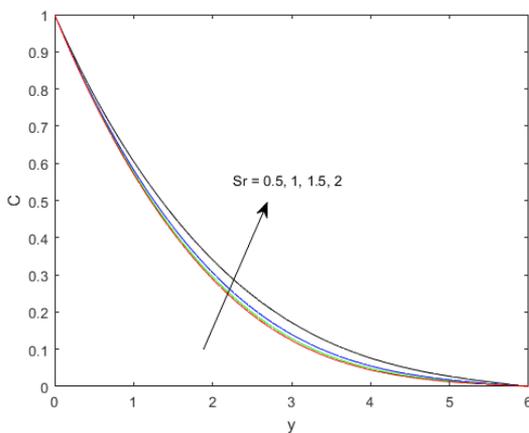


Figure 10. Concentration profiles for Sr

The wall velocity gradient which represent skin friction at the plate together with wall temperature gradient i.e. Nusselt number and gradient in the concentration in the vicinity of the plate, for various parameters of interest are depicted in Tables 2-4. In these tables, we notice that there is a reduction in the coefficient of skin friction with a rise in the values of thermal Grashof number, solutal Grashof number, Soret number, radiation parameter while it increases for magnetic parameter. Rate of heat transfer and rate of mass transfer increases for chemical reaction parameter and Schmidt number while these quantities decrease with radiation parameter and Soret number.

6. CONCLUSION

The current study focuses on the theoretical investigation of Soret effect on transient magnetohydrodynamic natural convection flow of an optically thick heat radiating, and chemically reactive nanofluid near a vertical infinite plate through a porous medium in the existence of second order homogeneous chemical reaction. The noteworthy findings of the boundary layer flow of the present study are briefed as follows:

- Throughout the boundary layer region, thermal and mass buoyancy forces, thermal radiation and thermo-diffusion effect tend to accelerate nanofluid velocity but completely opposite effect is observed for magnetic field and chemical reaction parameters.
- Tendency of thermal radiation is to enhance the fluid temperature in the boundary layer region.
- Mass diffusion and chemical reaction effects tend to reduce the species concentration of the nanofluid whereas a totally reverse effect is noticed on it for thermo-diffusion effect.
- We see a reduction in the Coefficient of Skin friction with a rise in the values of thermal Grashof number, solutal Grashof number, Soret number, radiation parameter while it increases for magnetic parameter. Rate of heat transfer and rate of mass transfer increases for chemical reaction parameter and Schmidt number while these quantities decrease with radiation parameter and Soret number.

Table 2. Skin friction coefficient (τ)

Gr	Gc	M	Kr	Sr	Nr	τ
2	4	3	3	0.5	4	1.2996
3						0.6132
5						0.2699
5	2	3	3	0.5	4	1.1546
	4					0.2699
	6					0.1257
5	4	1	3	0.5	4	0.1983
		3				0.2699
		5				0.3416
5	4	3	1	0.5	4	0.7867
			3			0.2699
			6			0.1239
5	4	3	3	0.2	4	0.5543
				0.5		0.2699
				0.8		0.1126
5	4	3	3	0.5	1	0.9657
					2	0.2699
					4	0.0219

Table 3. Rate of heat transfer (Nu)

Nr	Kr	- Nu
1	0.3	0.7845
2		0.4586
4		0.2451
4	1	0.1723
	3	0.3984
	6	0.5126

Table 4. Rate of mass transfer (S_h)

Sc	Kr	Sr	- S_h
0.22	3	0.5	0.1838
0.44			0.3746
0.66			0.5186
0.22	1	0.5	0.0914
	3		0.1838
	6		0.2435
0.22	3	0.2	0.4923
		0.5	0.1838
		0.8	0.0547

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NOMENCLATURE

M	Magnetic parameter
Sc	Schmidt number
g	Acceleration due to gravity
B_0	Uniform magnetic field
Pr	Prandtl number
Kr	Chemical reaction parameter
Nr	Radiation parameter
Sr	Soret number
Gr	Thermal Grashof number
Gm	Solutal Grashof number
k'	Permeability of porous medium
k^*	Rosseland mean absorption co-efficient
σ^*	Stefan-Boltzmann constant
u'	Fluid velocity
T'	Fluid Temperature
C'	Species concentration

Greek symbols

ρ_f	Density of base fluid
ρ_{nf}	Density of nanofluid
μ_{nf}	Dynamic viscosity of nanofluid
μ_f	Dynamic viscosity of base fluid
σ_{nf}	Electrical conductivity of nanofluid
σ_f	Electrical conductivity of base fluid
ψ	Nanoparticle volume fraction

Subscripts

s	Nanoparticle
f	Base fluid
nf	Nanofluid